

# High-precision studies in fundamental physics with slow neutrons

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ILL, 20 September 2016

# Topics

- The impossible particle and its properties
- Search for an electric dipole moment of the neutron
- Short-range gravity
- Test of Einstein's  $E = mc^2$

# The neutron before Chadwick



# The neutron before Chadwick



“Such an atom would possess striking properties. Its outer field would vanish [...] and therefore it should easily penetrate matter. The existence of such an atom is presumably difficult to observe with a spectrograph, and ...”

(„Nuclear Constitution of Atoms“, Proc. Royal Soc. 1920)

# The neutron before Chadwick



"Such an atom would possess striking properties. Its outer field would vanish [...] and therefore it should easily penetrate matter. The existence of such an atom is presumably difficult to observe with a spectrograph, and it could not be stored in a closed vessel."

(„Nuclear Constitution of Atoms“, Proc. Royal Soc. 1920)

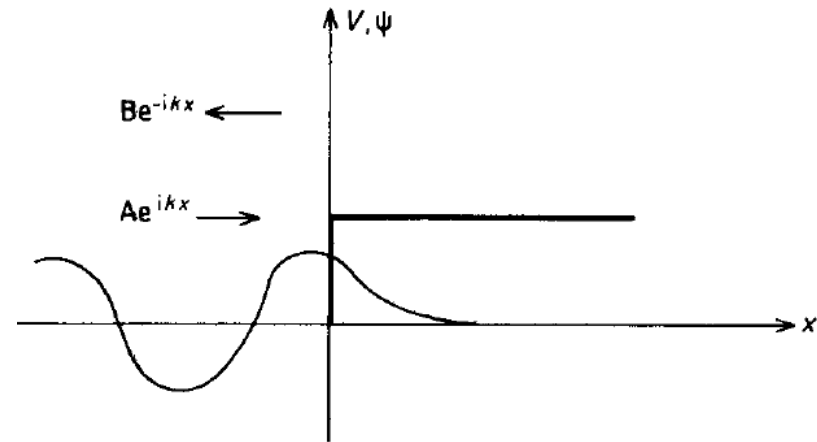
# How to store it nevertheless?

Mirror reflection under any angle of incidence

→ *UCN can be trapped in “neutron bottles”*

Trapping potential #1:

neutron optical potential  $V + iW$



Physical origin:

- neutron scattering by nuclei
- interference of incident and scattered waves
- refractive index:

$$n = \frac{k'}{k} = \sqrt{1 - \frac{V}{E}} \quad V = \frac{2\pi\hbar^2}{m} Na$$

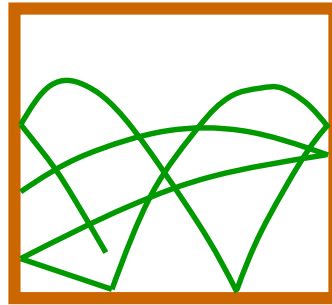
Typical values for V:

Be: 252 neV, Al: 54 neV, Ti: -49 neV

Trapping potential #2:

**neutron gravity** *mgz*

for  $\Delta z = 1$  m:  $\Delta E = 100$  neV



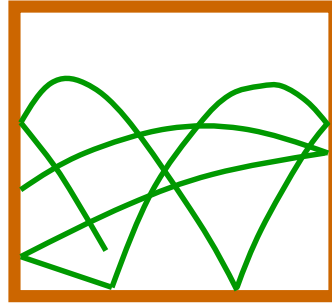
as good for trapping  
(if bottle is tall enough):



Trapping potential #2:

**neutron gravity** *mgz*

for  $\Delta z = 1$  m:  $\Delta E = 100$  neV





Trapping potential #3:

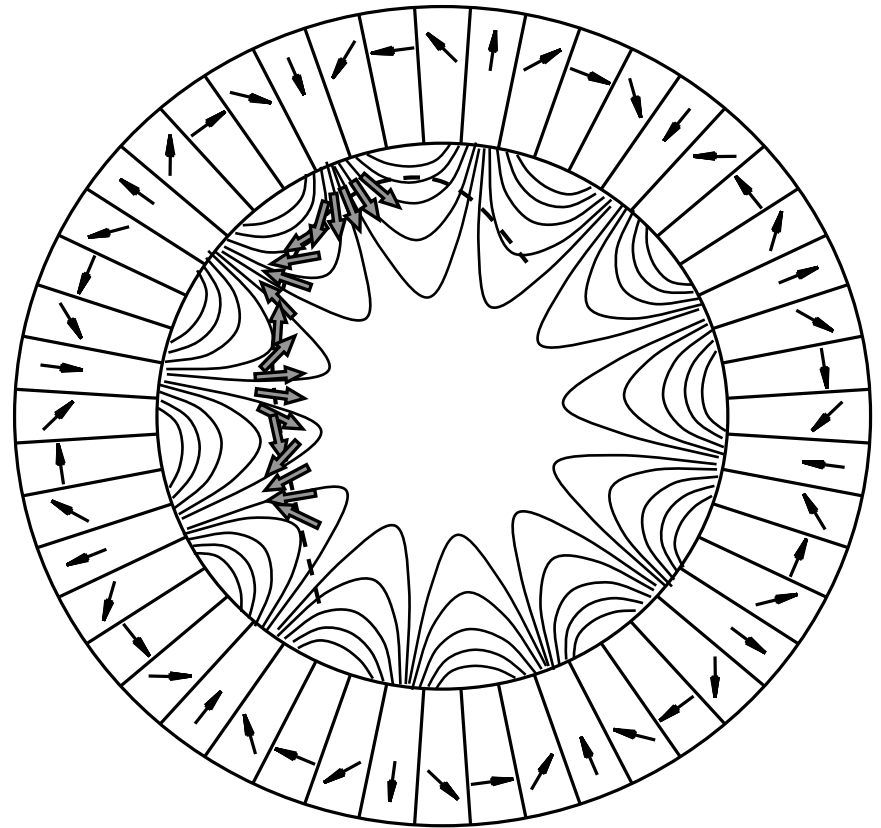
**magnetic interaction**  $\pm\mu B$

for  $\Delta B = 1 \text{ T}$ :  $\Delta E = \pm 60 \text{ neV}$

Adiabatic spin transport if

$$\frac{1}{|B|} \cdot \left| \frac{dB}{dt} \right| \ll \frac{\mu \cdot B}{\hbar} = \omega_L$$

→ mT fields sufficient in typical situations



Magnetic gradient fields suppress losses due to wall collisions

# Neutron properties:

Property	Symbol	Value
Spin <sup>Parity</sup>	$s^P$	$\frac{1}{2}^+$
Mass (relative to $^{12}\text{C}$ mass standard)	$m_n$	1.008 664 915 8(6) u
Mass (absolute units)		939.565 33(4) MeV $c^{-2}$
Neutron - proton mass difference	$m_n - m_p$	0.001 388 448 9(6) u 1.293 331 8(5) MeV $c^{-2}$
Charge	$q_n$	$(-0.4 \pm 1.1) \times 10^{-21} e$
Mean-square charge radius	$\langle r_n^2 \rangle$	-0.116 1(22) fm <sup>2</sup>
Electric polarisability	$\alpha_n$	$(9.8_{-2.3}^{+1.9}) \times 10^{-4} \text{ fm}^3$
Magnetic moment	$\mu_n$	-1.913 042 7(5) $\mu_N$ $= -6.030 773 8(15) \times 10^{-8} \text{ eV T}^{-1}$
Electric dipole moment	$d_n$	$< 2.9 \times 10^{-26} \text{ e cm (90\% c.l.)}$
Mean $n\bar{n}$ -oscillation time of free neutron	$\tau_{n\bar{n}}$	$> 8.6 \times 10^7 \text{ s (90\% c.l.)}$
... of bound neutron		$> 1.2 \times 10^8 \text{ s (90\% c.l.)}$
Parameters of $\beta$ -decay, $n \rightarrow p + e^- + \bar{\nu}_e$		
$Q$ -value	$Q$	0.782 332 9(5) MeV $c^{-2}$
Mean life time	$\tau_n$	885.7(8) s
Ratio of weak coupling constants $g_A/g_V$	$\lambda$	-1.2670 (30)
Coefficients of angular correlations:		
neutron spin - electron momentum: $P_n \cdot p_e$	$A$	-0.1162 (13)
momenta of antineutrino and electron: $p_\nu \cdot p_e$	$a$	-0.102 (5)
neutron spin - antineutrino momentum	$B$	0.983 (4)
triple correlation $P_n \cdot (p_e \times p_\nu)$	$D$	$-0.6 (10) \times 10^{-3}$
Phase angle between $V$ and $A$ weak currents	$\phi_{VA}$	$-180.08 (10)^0$

# The Big Bang

15 thousand million years

A world of matter

1 thousand million years

300 thousand years

3 minutes

1 second

**n̄ oscillations**

**neutron lifetime**

**nEDM**  $10^{-10}$  seconds

$10^{-34}$  seconds

$10^{-43}$  seconds

$10^{32}$  degrees

$10^{27}$  degrees

$10^{15}$  degrees

$10^{10}$  degrees

$10^9$  degrees

6000 degrees

**nuclear few-body interactions**

18 degrees

**Heavy elements**

???

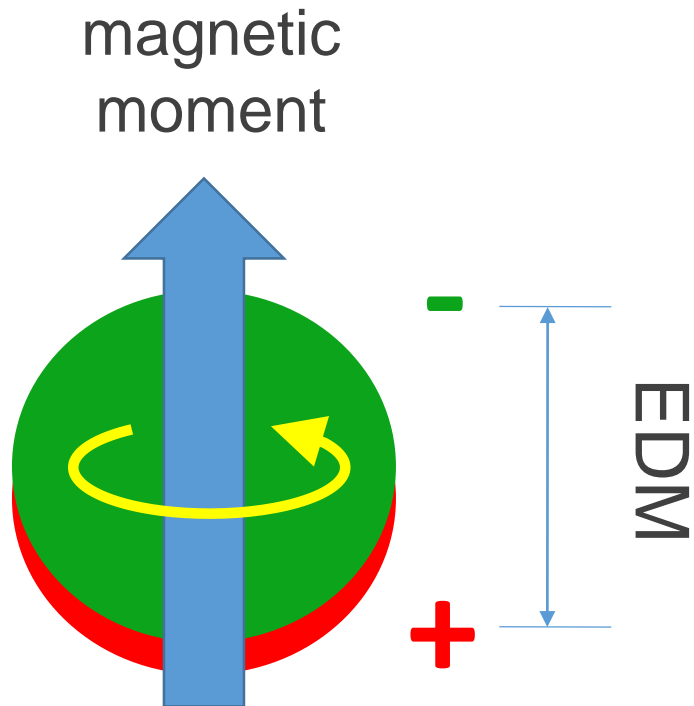
3 degrees K

- radiation
- particles
- $W^+$  } heavy particles carrying the weak force
- $W^-$  }
- $Z$  }
- quark
- anti-quark
- $e^-$  electron

- $\bar{e}$  positron (anti-electron)
- proton
- neutron
- meson
- H hydrogen
- D deuterium
- He helium
- Li lithium

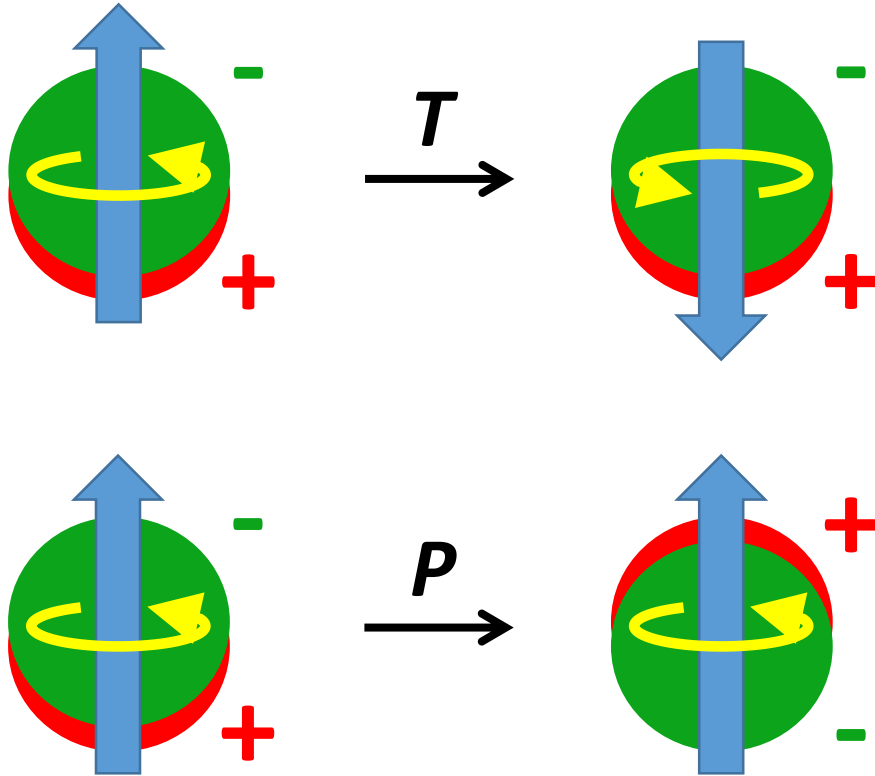
M. S. Turner

# Search for an electric dipole moment of the neutron



# Violation of fundamental symmetries

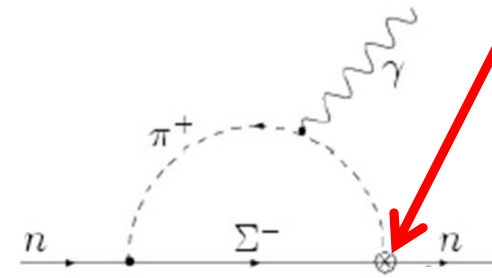
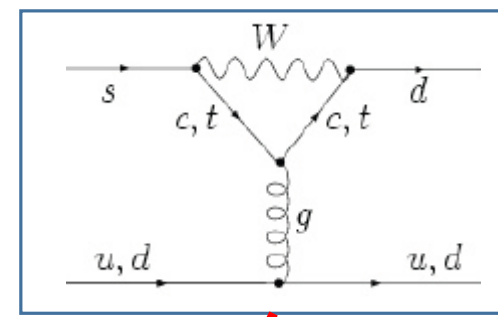
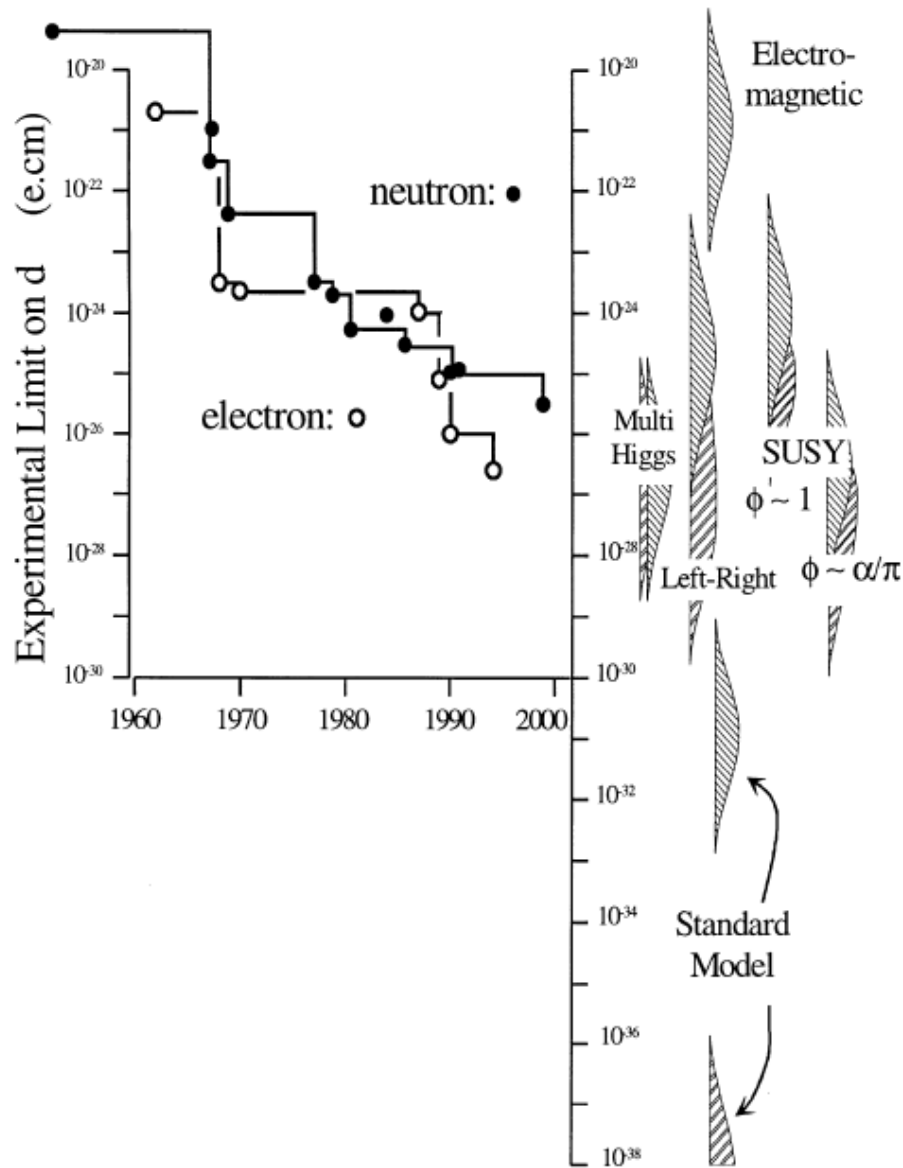
Purcell and Ramsey, PR 78 (1950) 807



- A non-zero particle EDM violates  $T$  (time reversal symmetry) and parity  $P$
- If we assume  $CPT$  conservation, also  $CP$  is violated, which is needed to explain the matter/antimatter asymmetry in the Universe

$$H = -\mu\mathbf{B} \cdot \frac{\mathbf{S}}{S} - d\mathbf{E} \cdot \frac{\mathbf{S}}{S}$$





- $CP$  violation within the Standard Model (SM) is too weak to explain the matter/antimatter asymmetry in the Universe
- nEDM tiny in the SM ( $10^{-31}$  ecm), but large in many beyond-SM theories
- nEDM sensitive probe to search new fundamental forces

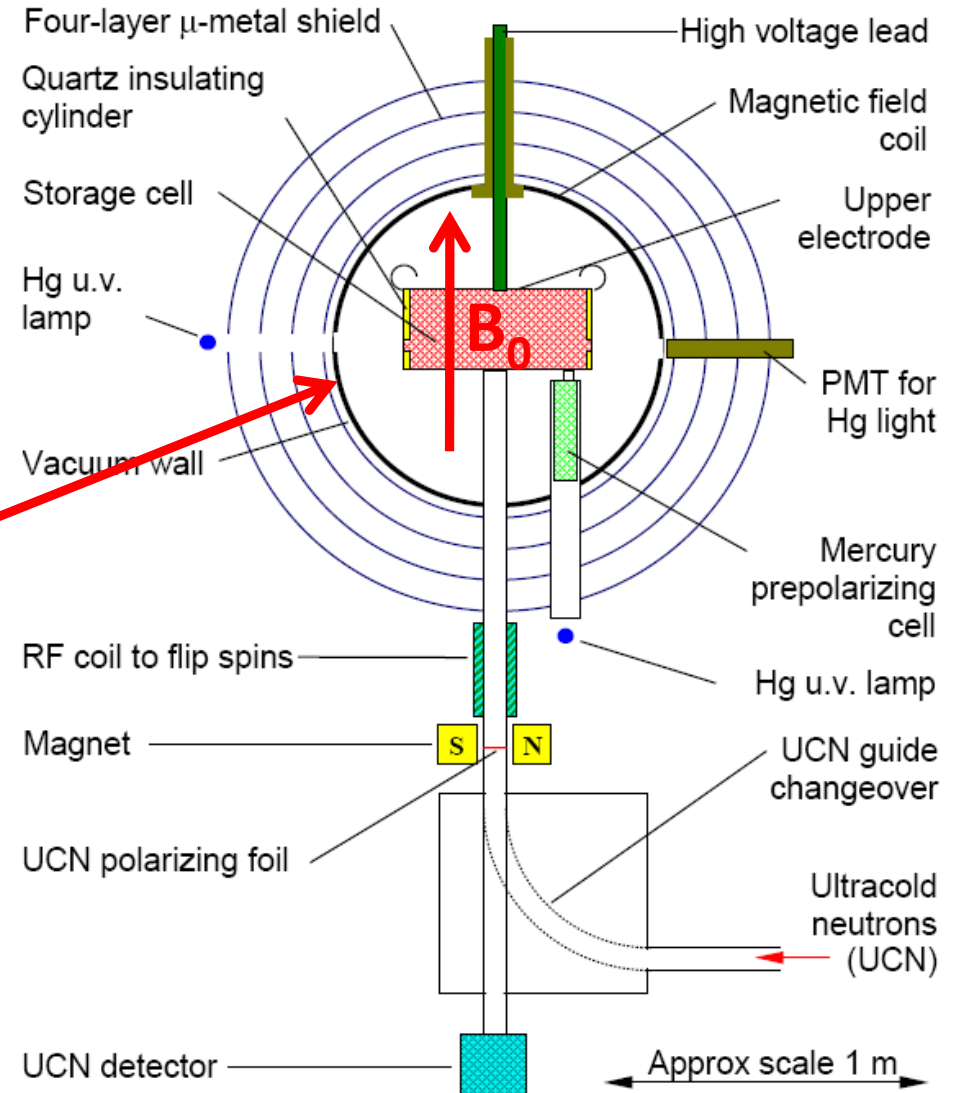
# How is it measured?

Ultra-cold neutrons (UCN)  
trapped at 300 K in vacuum



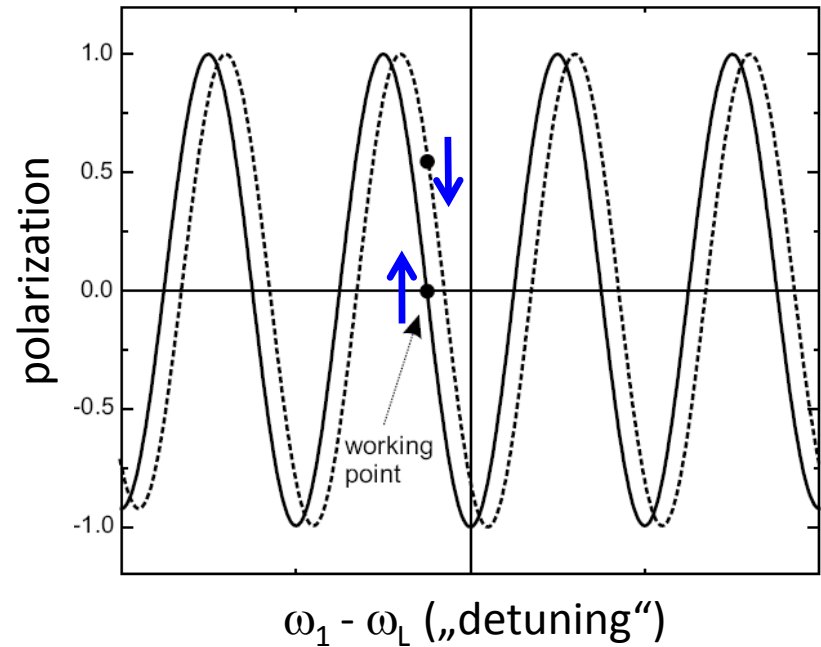
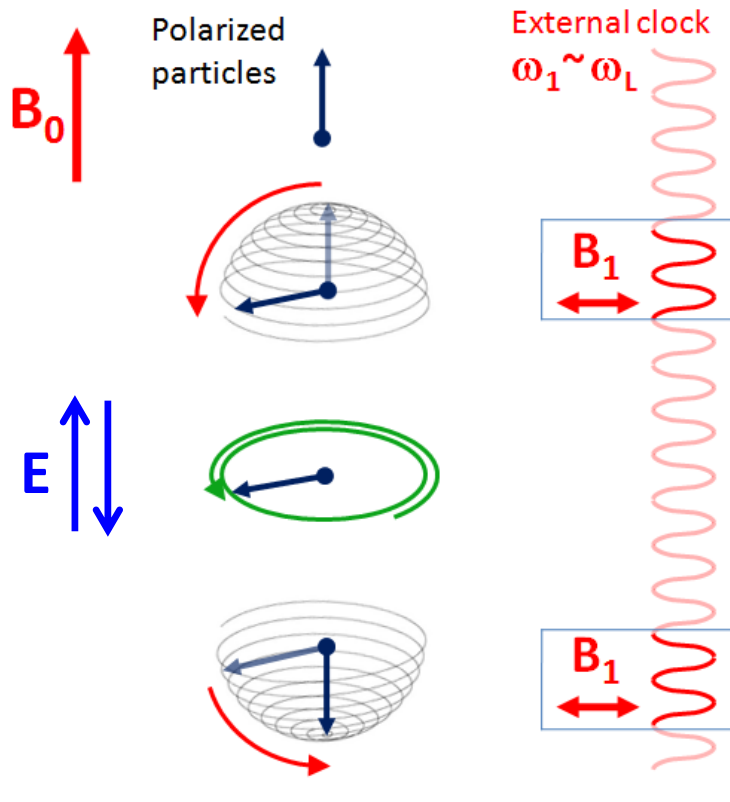
~ 0.5 m  
↔

RAL/SUSSEX/ILL experiment:



# Ramsey's method

Particle beam or trapped particles  
(...spin echo)



Experimental sensitivity:

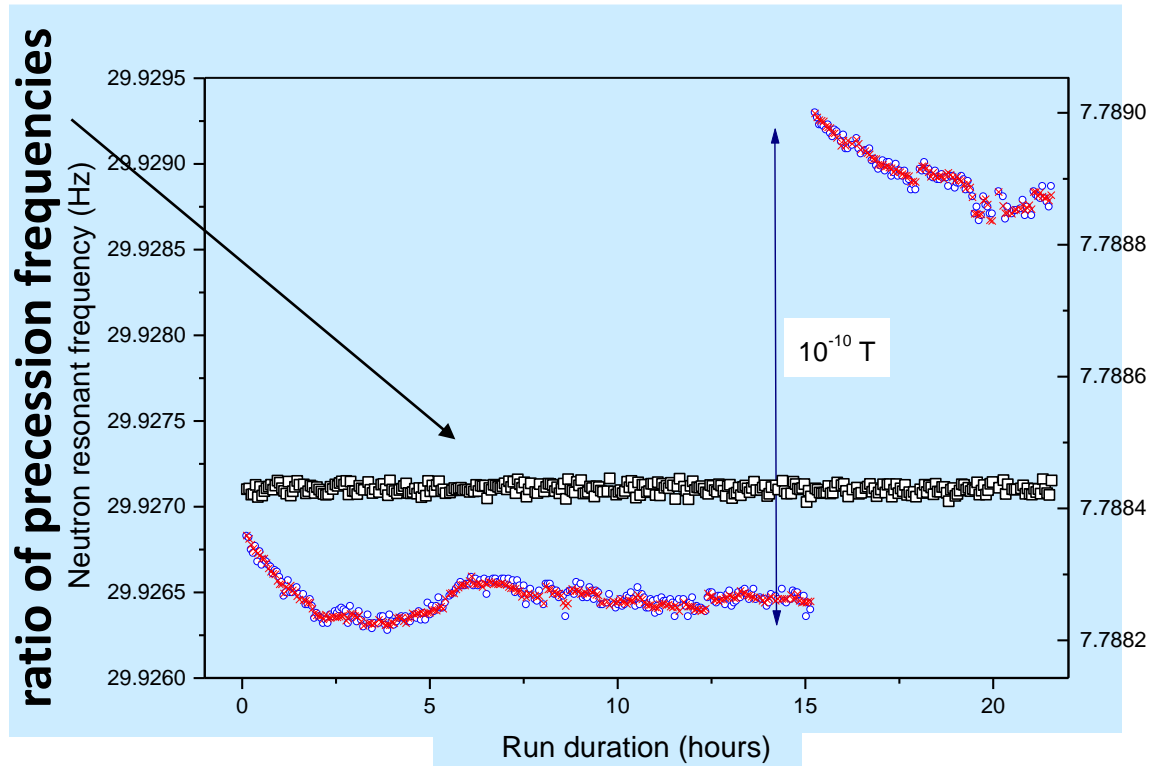
$$\sigma_{d_n} = \frac{\hbar}{2\alpha ET \sqrt{N}}$$

EDM changes frequency:

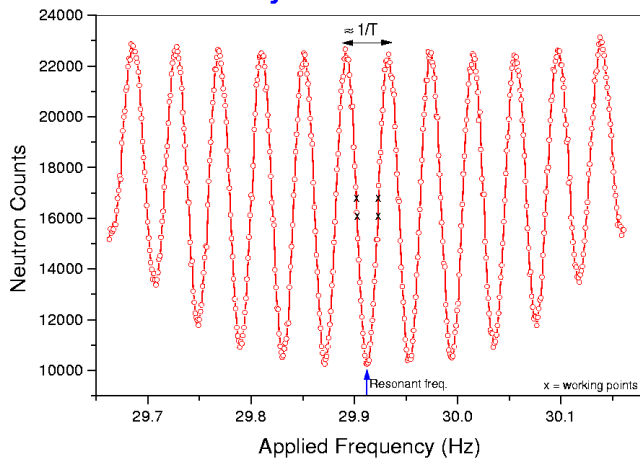
$$\hbar\omega_L \sim \mu_n B \pm d_n E$$



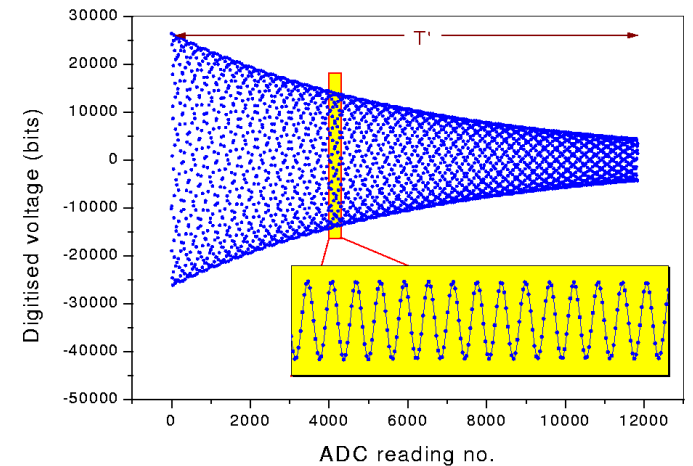
# $^{199}\text{Hg}$ co-magnetometer for correction of magnetic field drifts



### Ramsey Resonance Curve



### Mercury spin precession

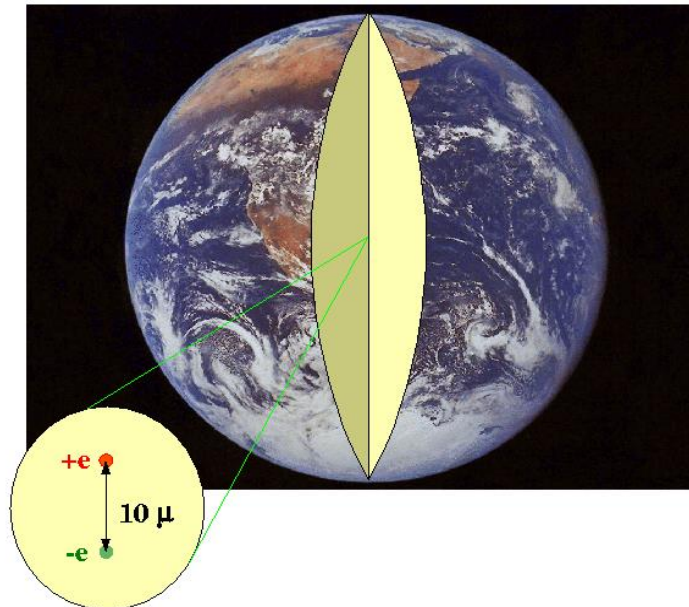


# Best result so far (RAL / Sussex / ILL)

$$|d_n| < 2.9 \times 10^{-26} \text{ e cm (90\% CL)}$$

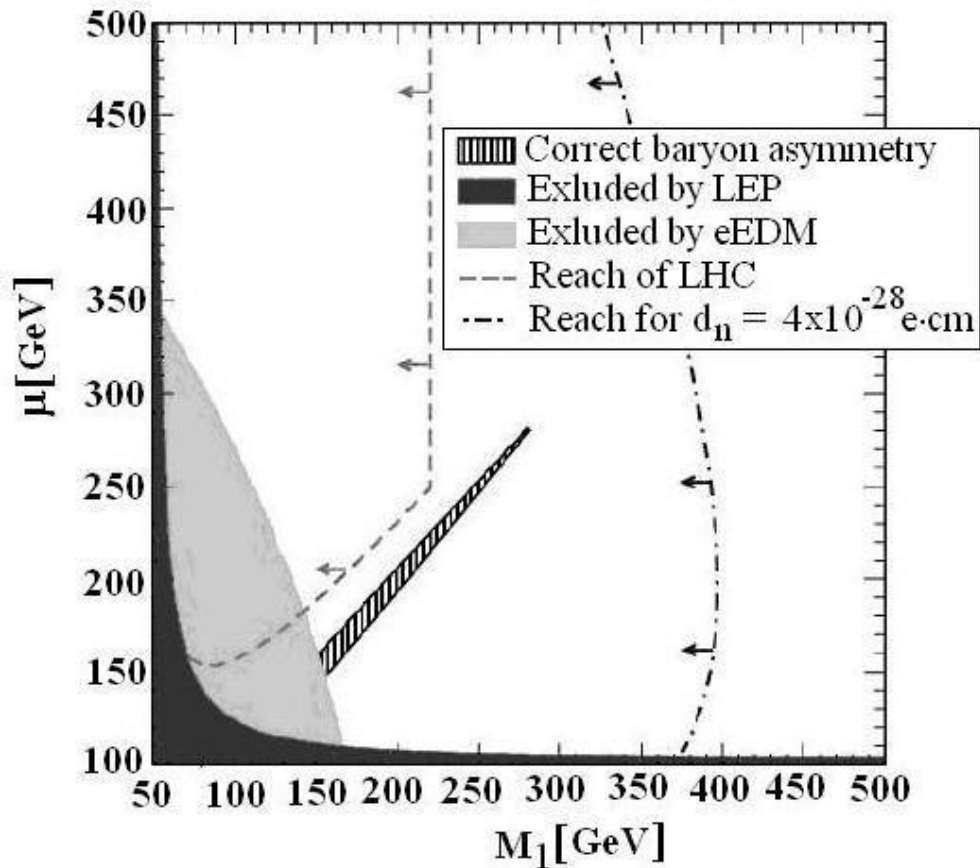
C.A. Baker et al., [PRL 63 \(2006\) 131801](#)

- $10^{-22}$  eV spin-dependent interaction
- one spin precession per half year



# Next steps?

- World-wide effort: projects at PSI, SNS, TRIUMF, TUM, PNPI, ILL
- Accuracy goal: below  $10^{-27}$  ecm
- needs new UCN sources, excellent magnetic shielding...



# Short-range gravity

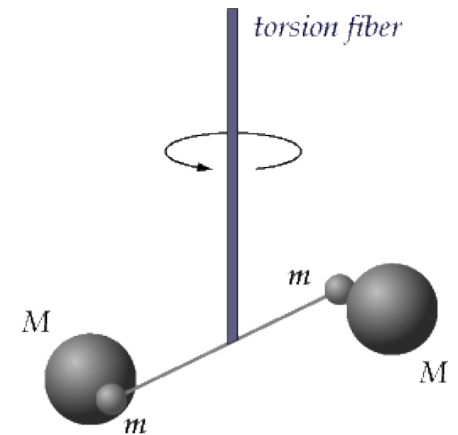
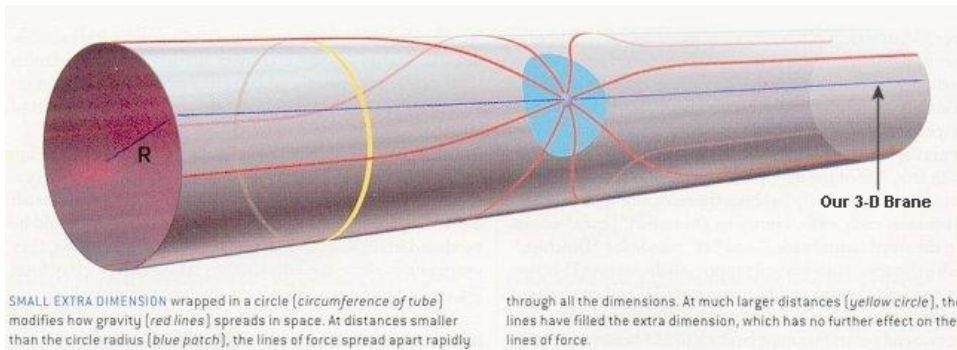
Small extra-dimensions:

Explanation why gravity is such a weak force?

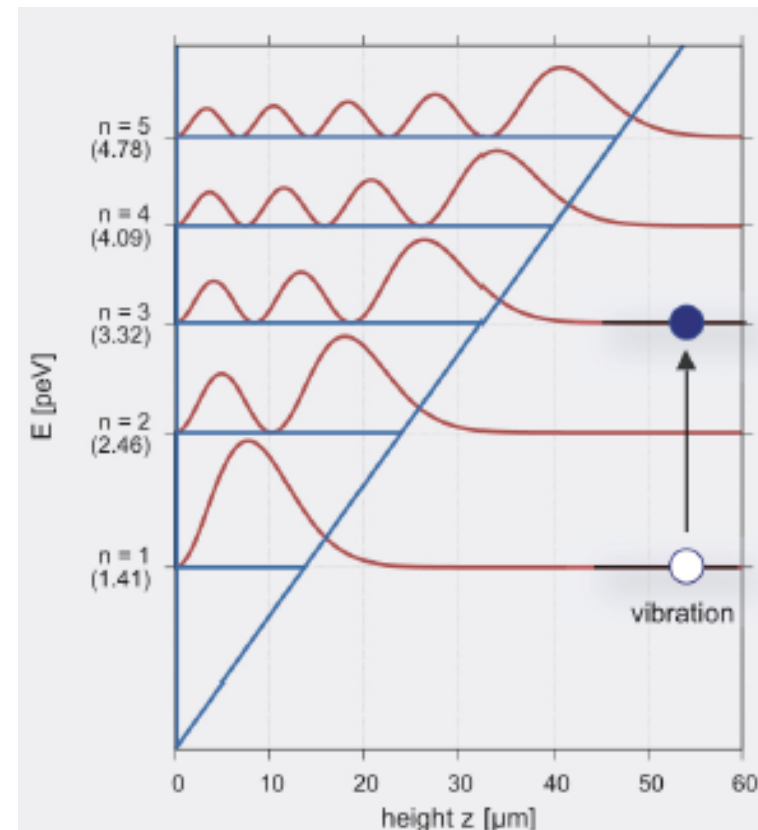
Modification of gravity

with  $n$  additional dimensions at distances  $r < R$ :

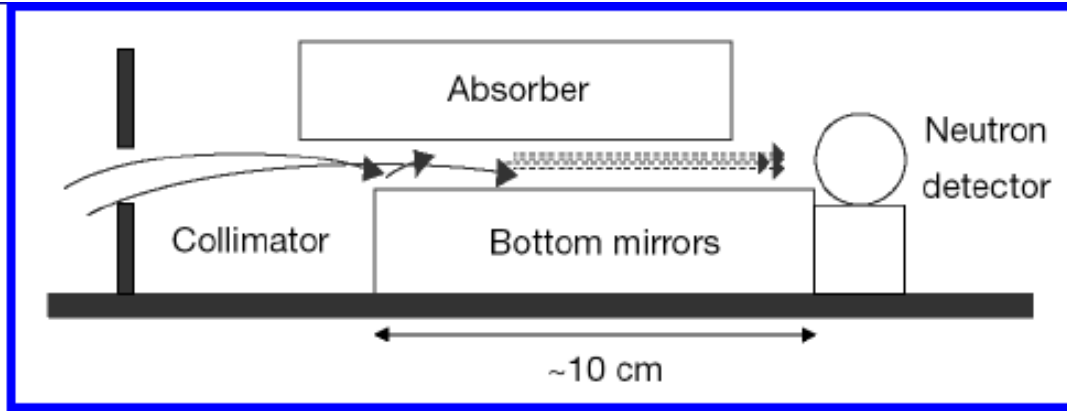
$$F = -G \frac{m_1 m_2}{r^2} \rightarrow -g \frac{m_1 m_2}{r^2} \frac{L^n}{r^n}$$



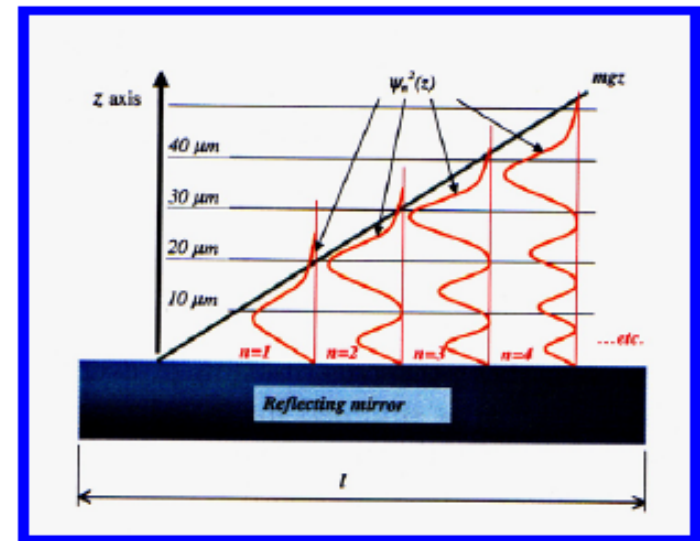
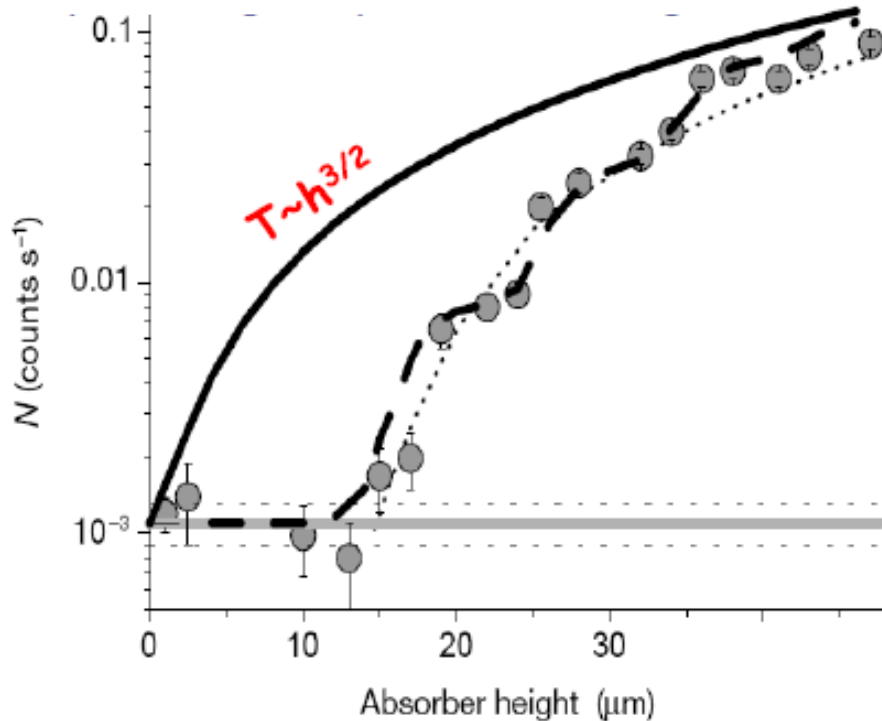
New spectroscopic tool:



# First observations (2002)



V. Nesvizhevsky et al., *Nature* **415** (2002) 299



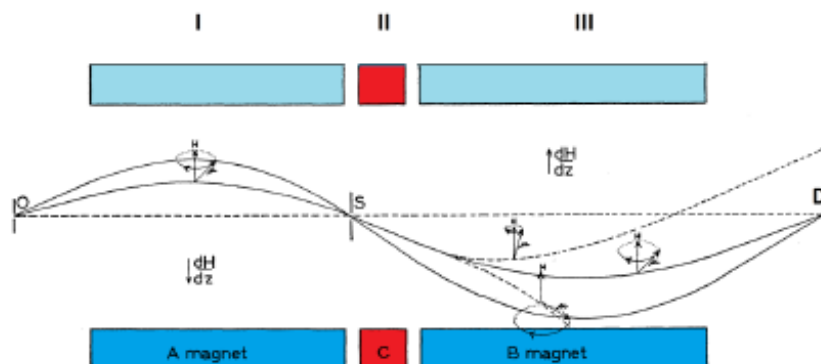


# Rabi-type spectroscopy of gravity

*q*Bounce collaboration (H. Abele, T. Jenke...)



## NMR Spectroscopy Technique to explore magnetic moments



### 3 Regions:

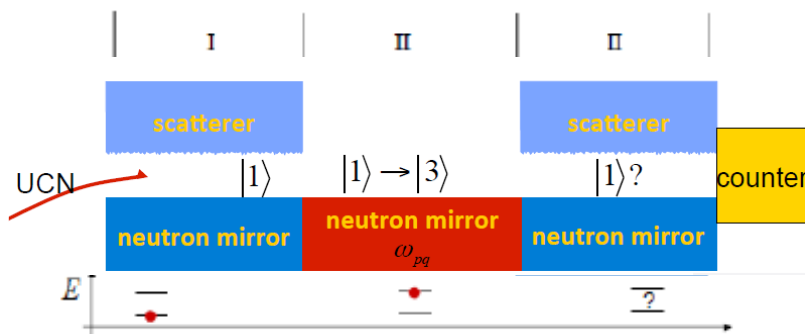
I: 1st State selector/ Polarizer

II: Coupling

– RF field

III: 2nd State Selector / Analyzer

## Gravity Resonance Spectroscopy Technique to explore gravity



### 3 Regions:

I: 1st State selector/ Polarizer

II: Coupling

– Vibr. mirror

— III: 2nd State Selector / Analyzer

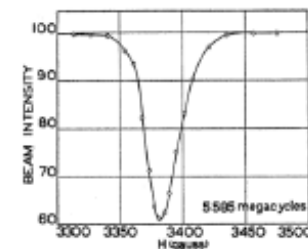


FIG. 4. Resonance curve of the  $\text{Li}^7$  nucleus observed in  $\text{LiCl}$ .

# Most recent results on Gravity Resonance Spectroscopy

Transitions 1-3 and 1-4 observed

1-3:  $(46 \pm 5)\%$  Intensity drop

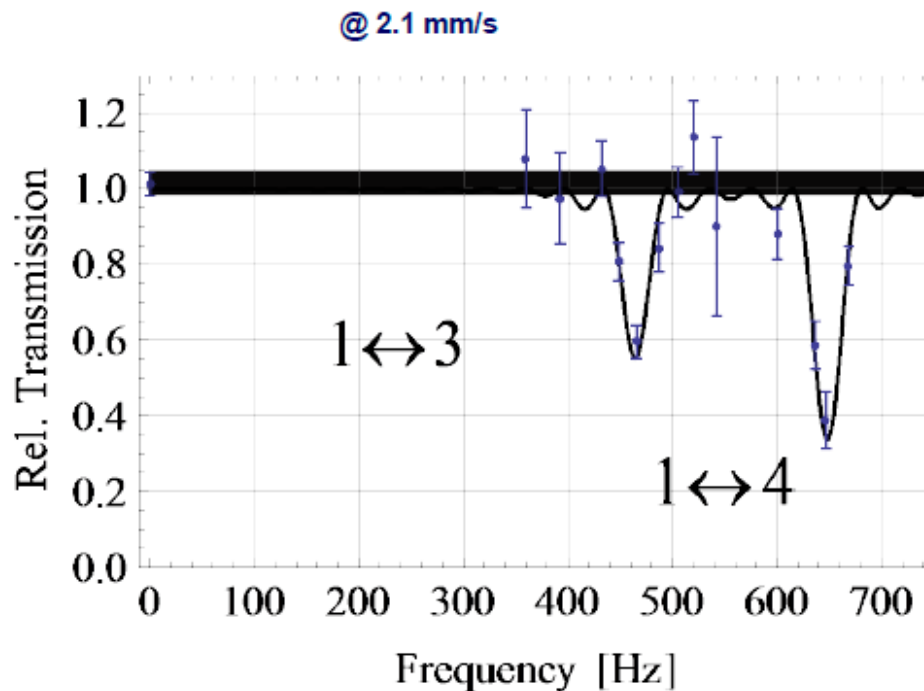
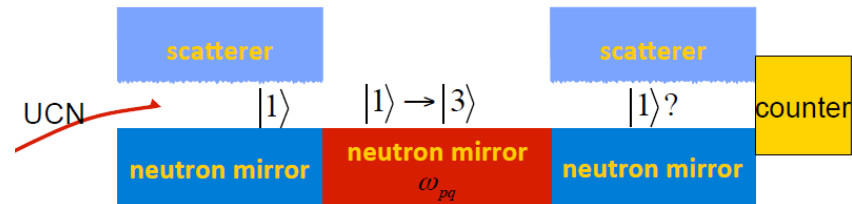
1-4:  $(61 \pm 7)\%$

60 measurements

Preliminary,

$$\nu_{13} = 463.74^{+1.05}_{-1.10} \text{ Hz}$$

$$\nu_{14} = 648.24^{+1.46}_{-1.53} \text{ Hz}$$



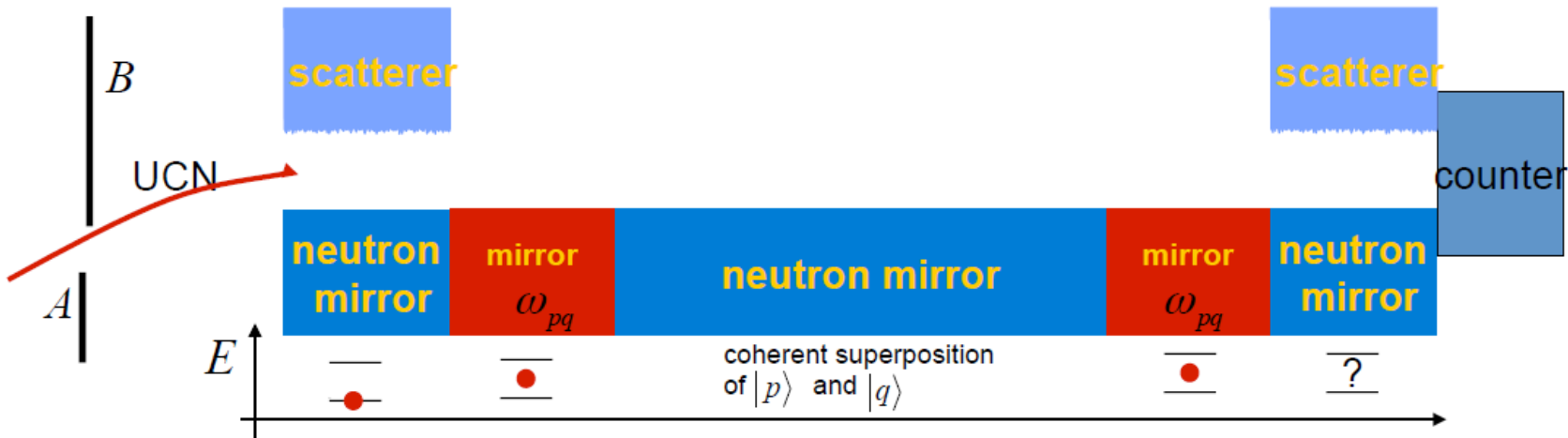
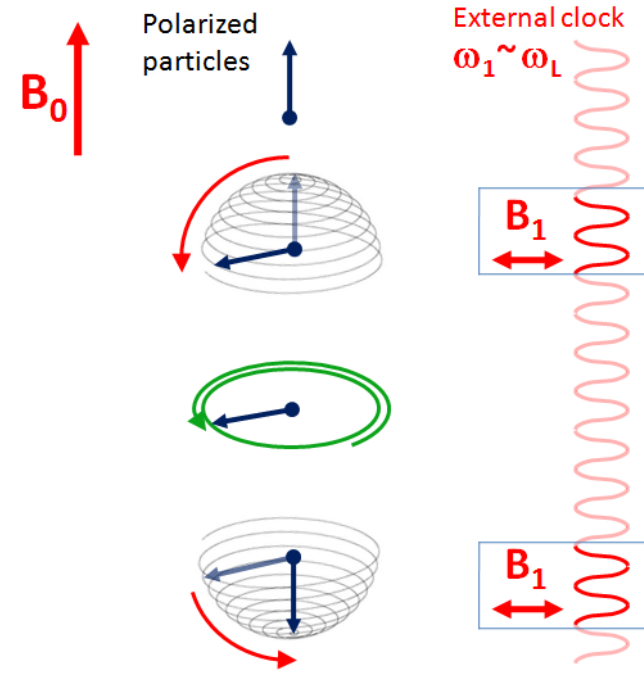
# Ramsey spectrometer for gravity states

H. Abele et al., *Phys. Rev. D* **81** (2010) 065019



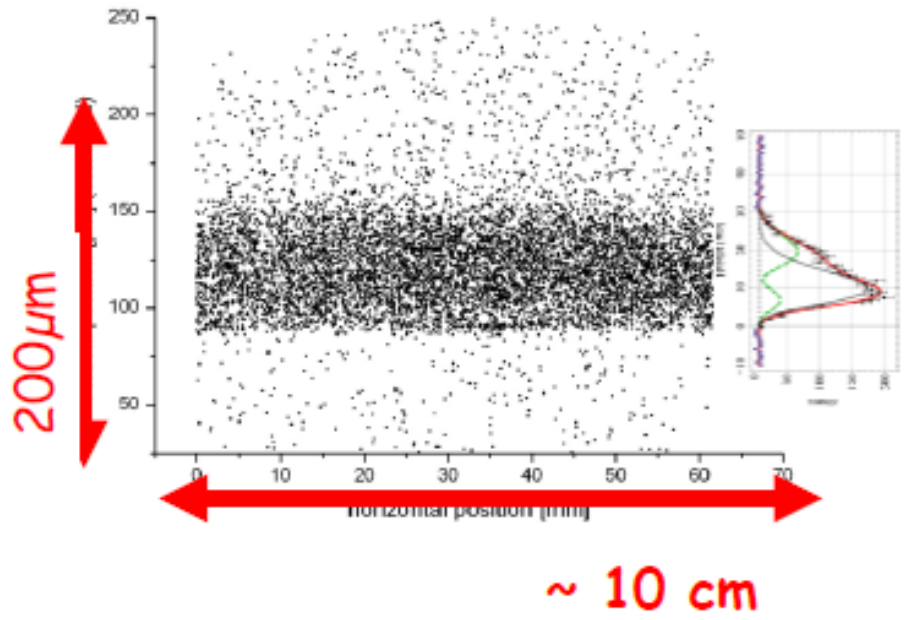
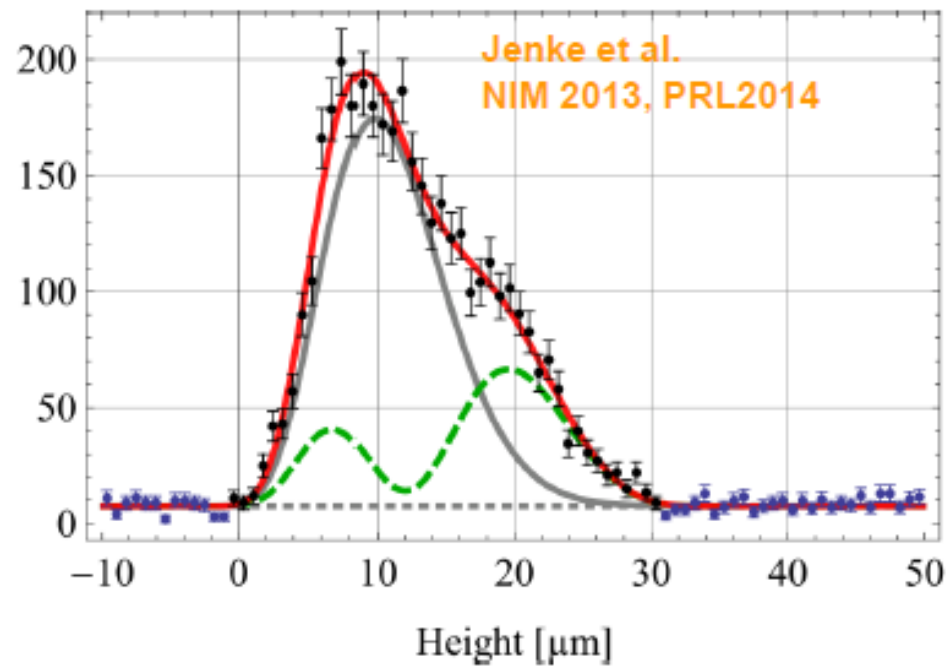
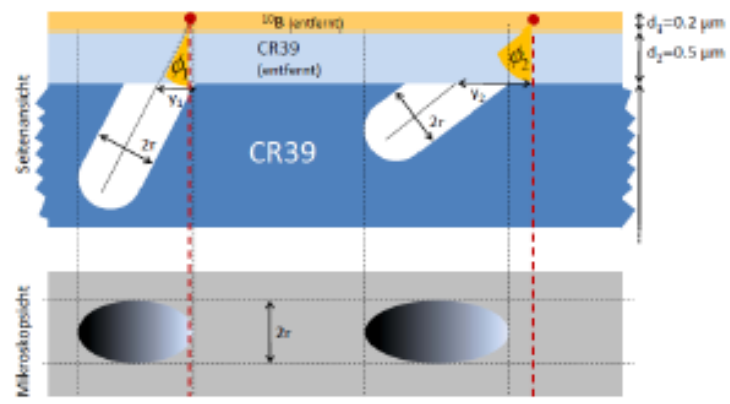
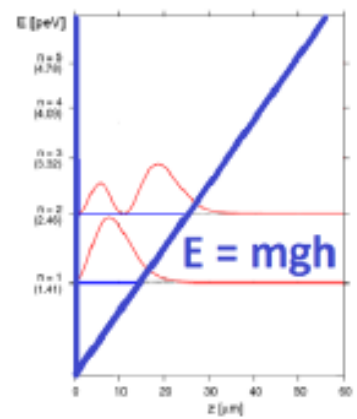
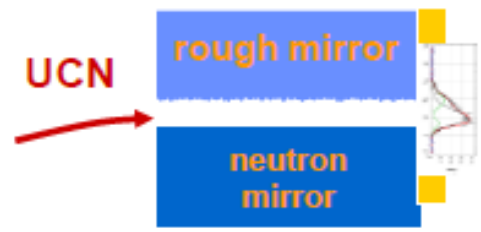
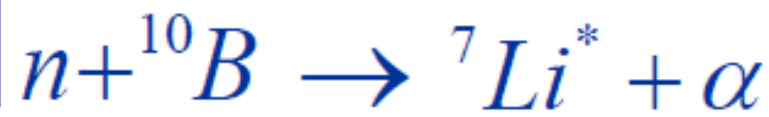
Advantages:

- Long flight path  $\rightarrow$  smaller uncertainty  $\Delta E$
- static central mirror for free state evolution





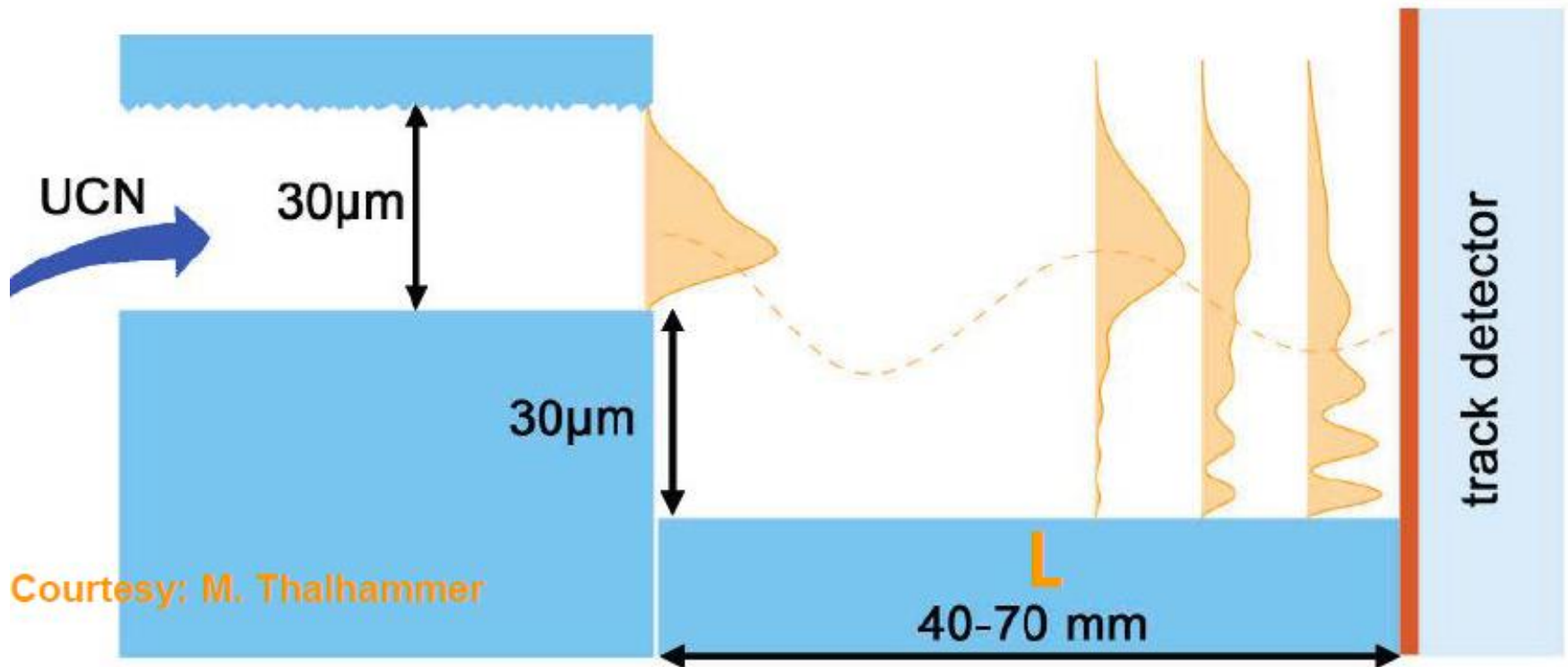
# Airy - Quantum States 1 & 2



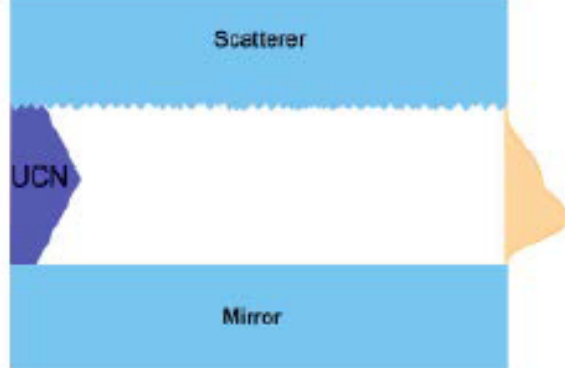
# Snapshots of $|\psi|^2$ with 1.5 $\mu\text{m}$ resolution

$$\Psi(z, t) = \sum_{n=0}^{\infty} c_n e^{-iE_n t/\hbar} \psi_n(z)$$

$$\psi_n(z) \sim A i \left[ \frac{z}{z_0} - \frac{E_n}{E_0} \right]; c_n = \int_0^{\infty} \Psi(z, 0) \psi(z) dz$$



# Preparation $L = 0$



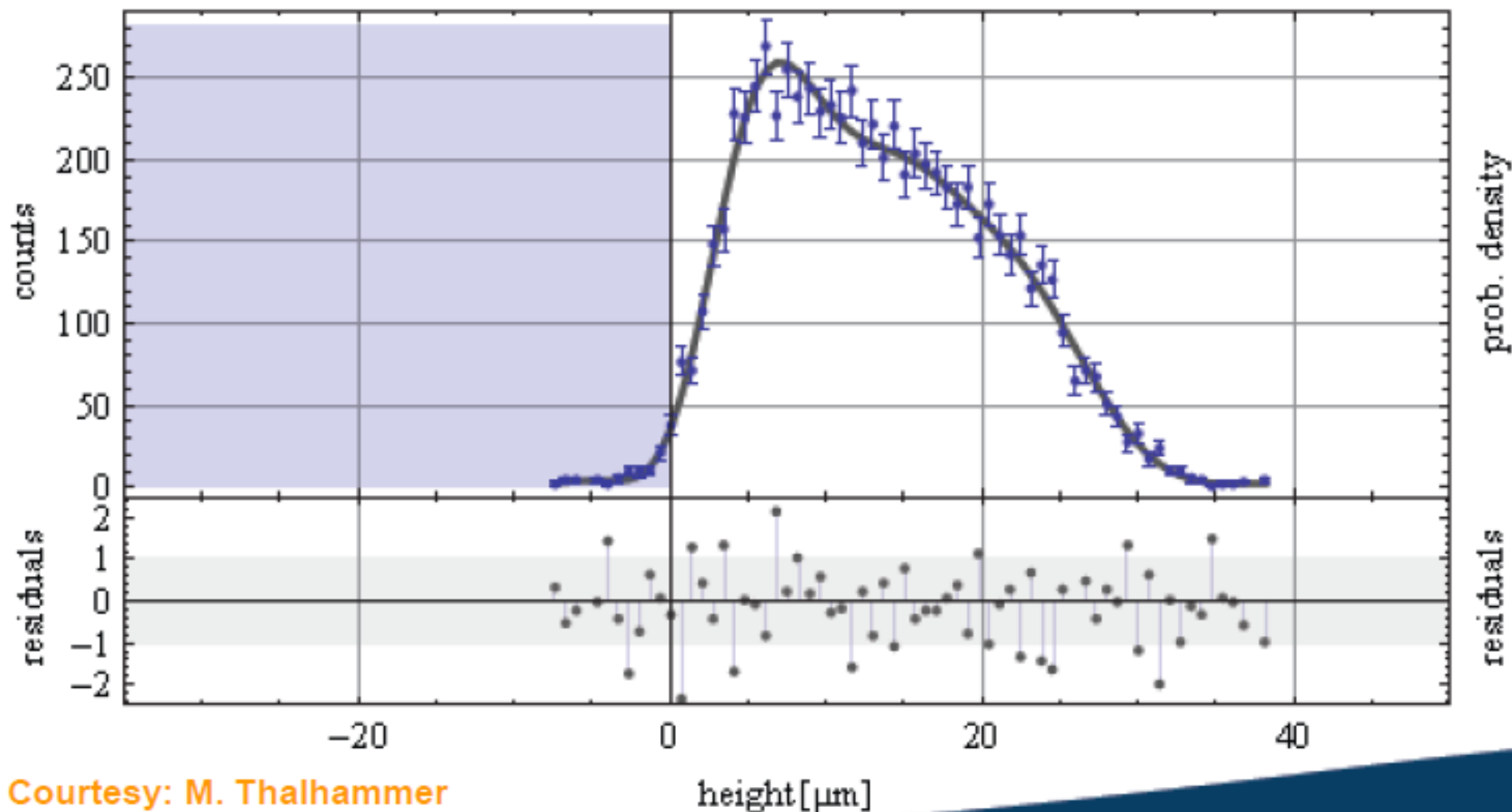
$$|\Psi_I(z, t_1)|^2 = \sum_n |c_n(t_1)|^2 \cdot |\psi_n(z)|^2$$

$$|c_1|^2 = 45\%$$

$$|c_2|^2 = 36\%$$

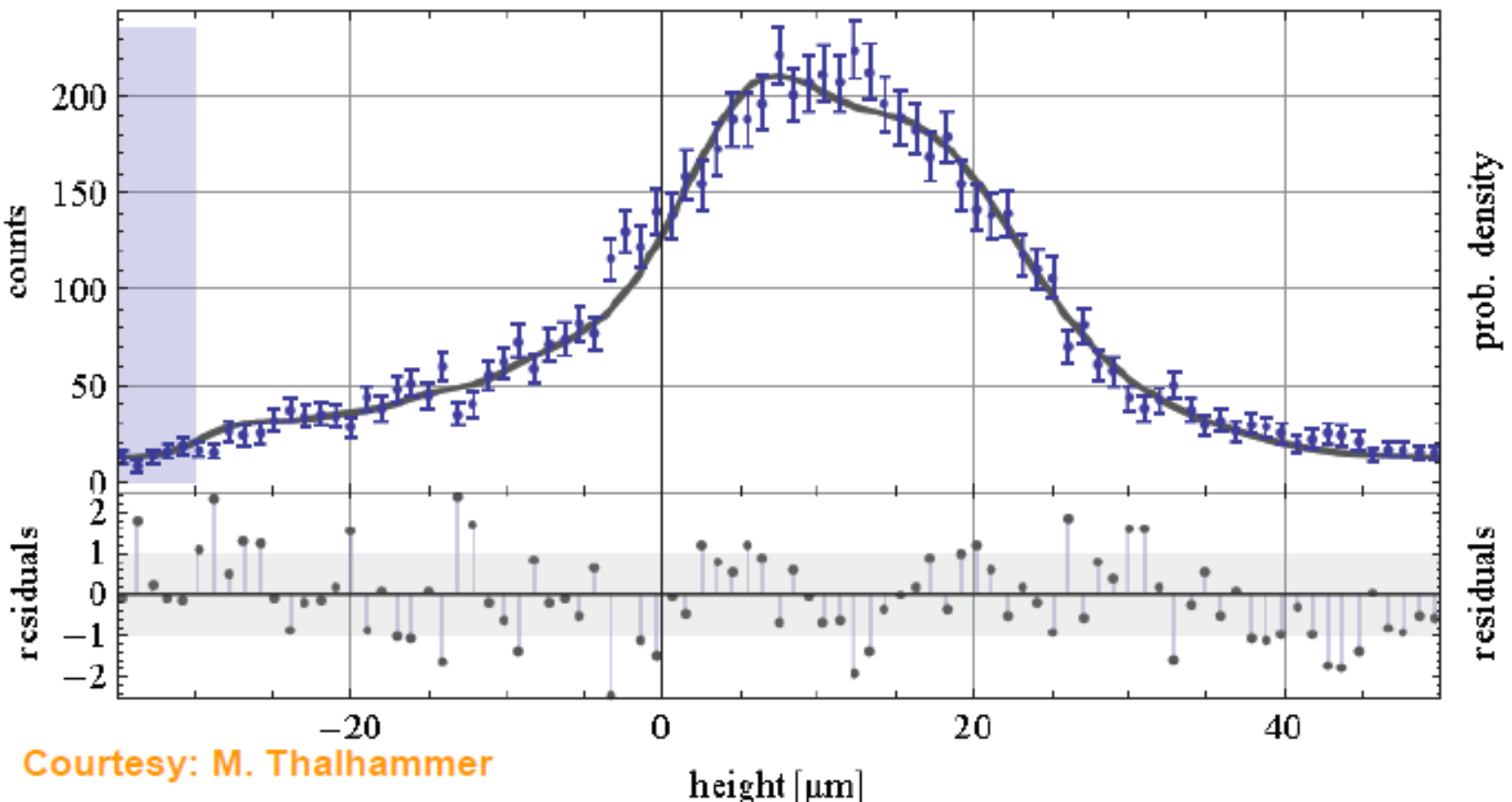
$$|c_3|^2 = 18\%$$

preliminary



Courtesy: M. Thalhammer

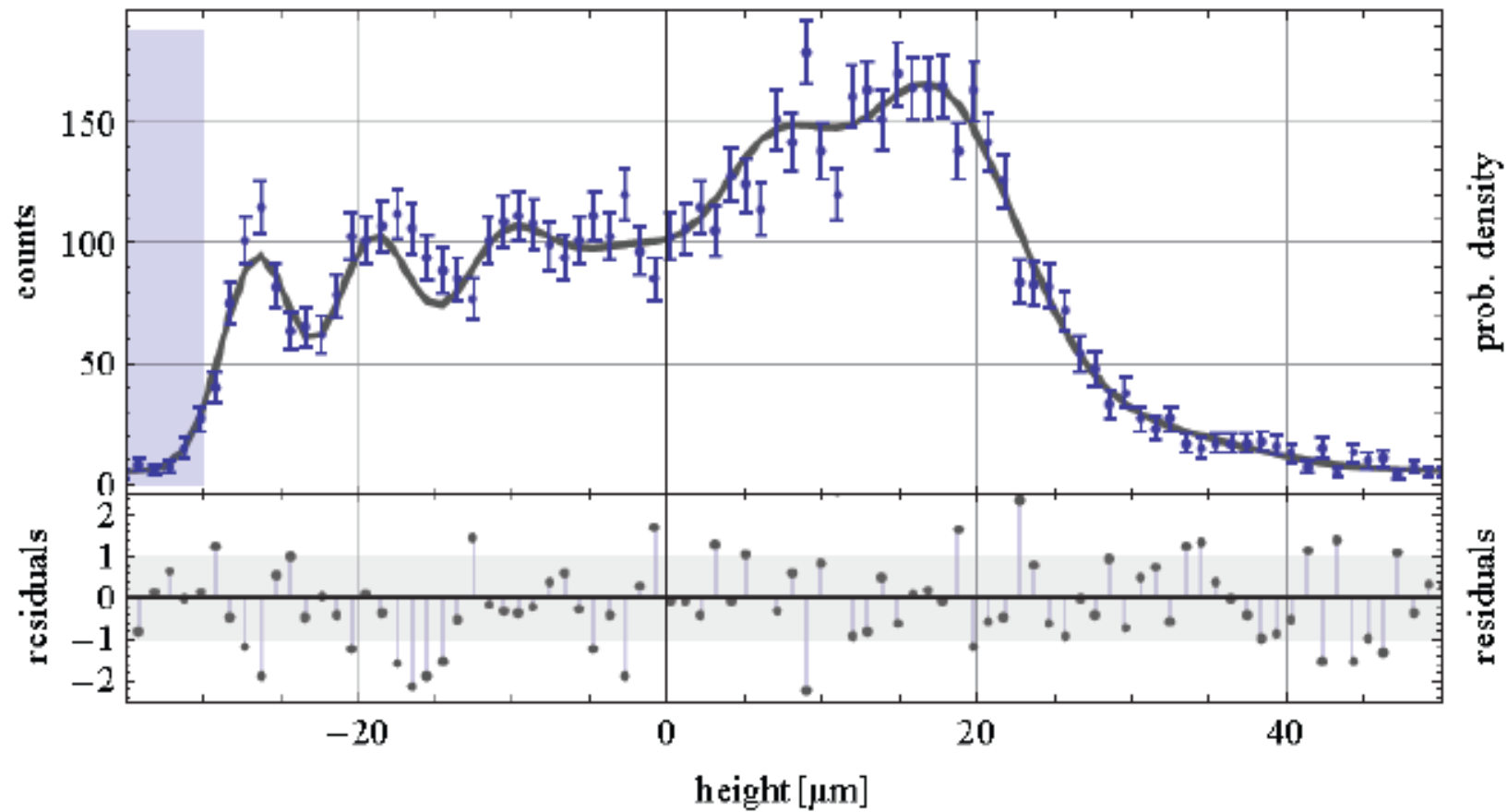
# 2nd bounce, 2nd turning point, $L = 41$ mm



Courtesy: M. Thalhammer

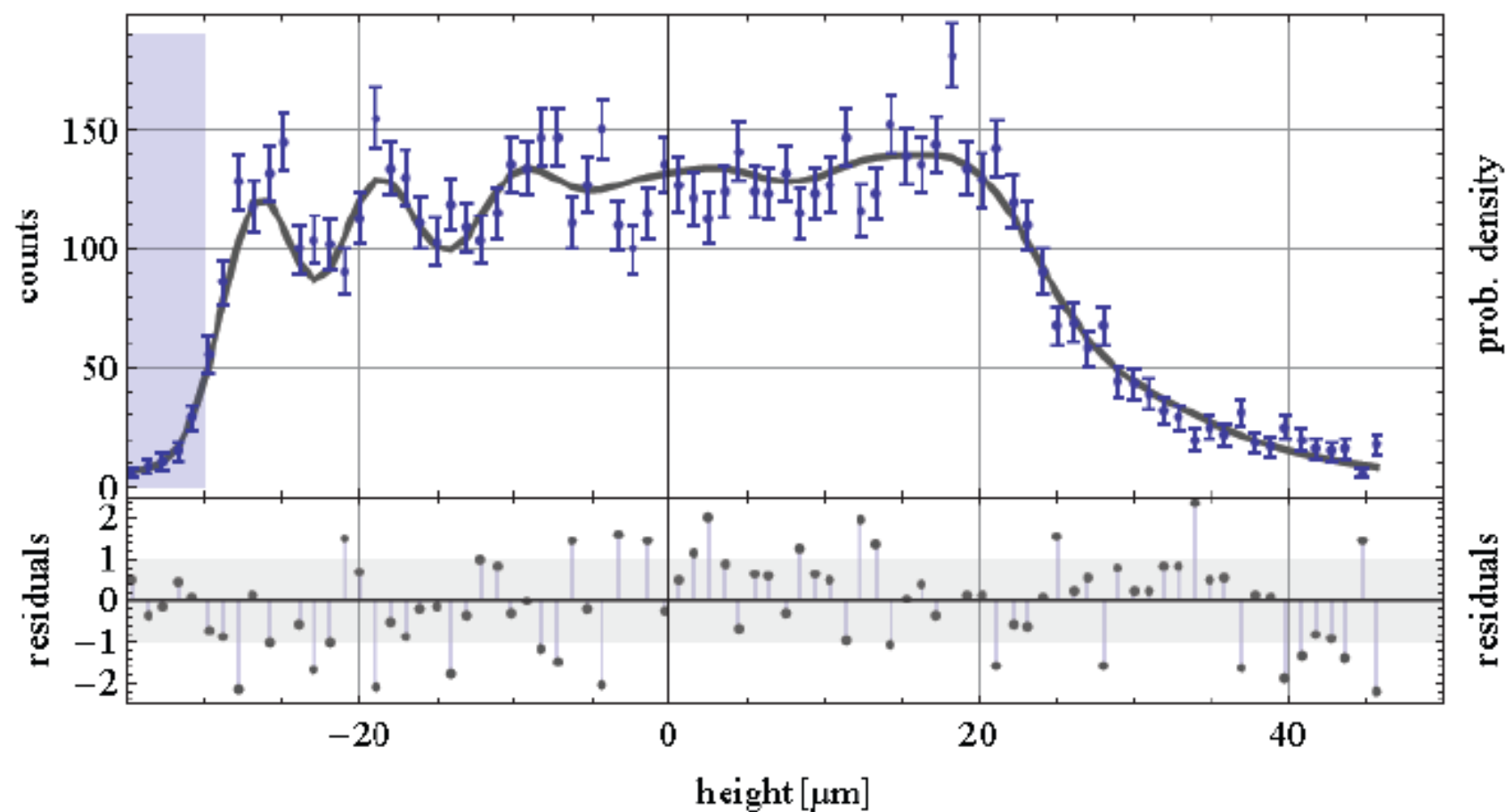
# Move downwards, $L = 51$ mm

51 mm

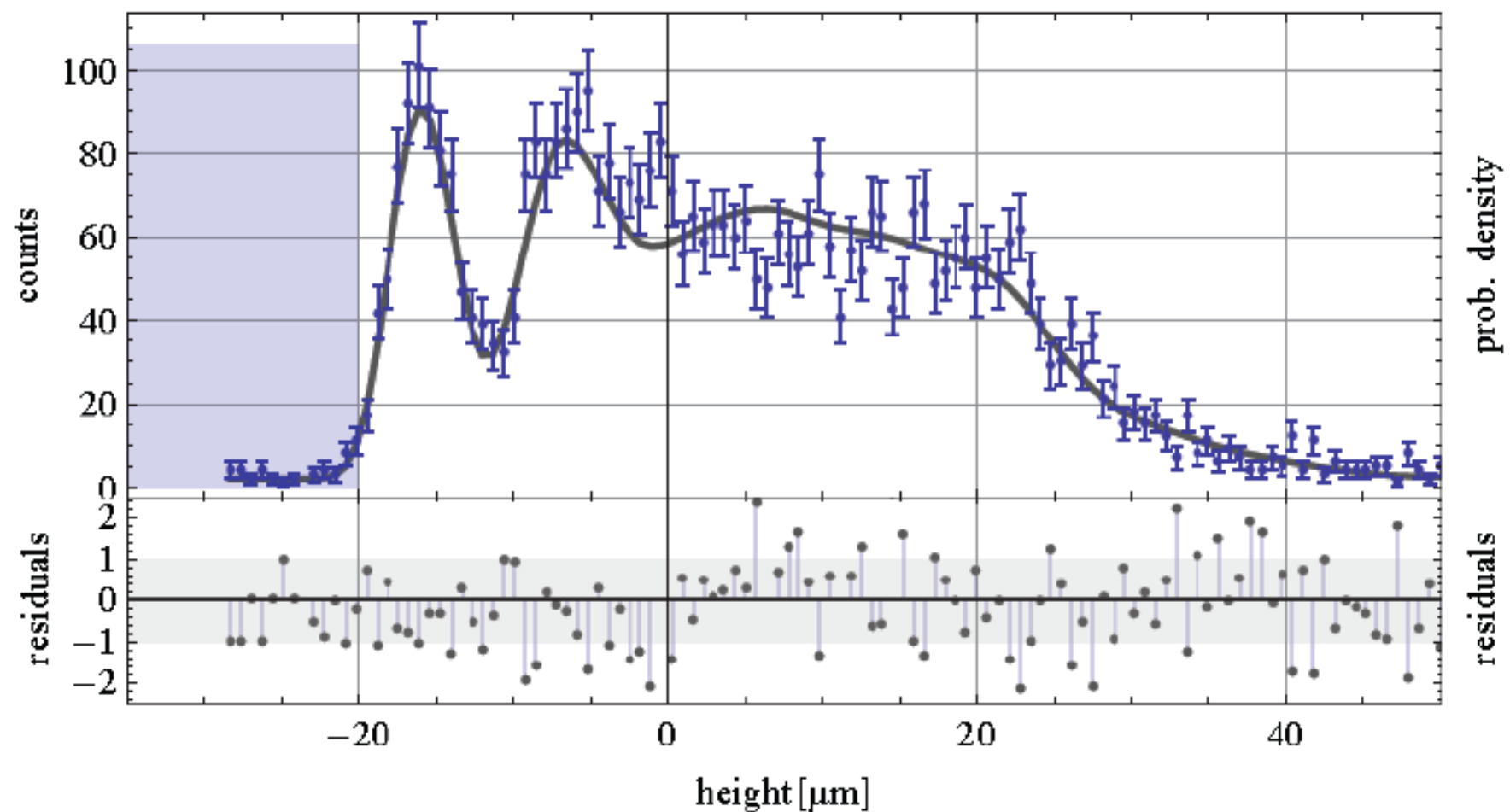


**L = 54 mm**

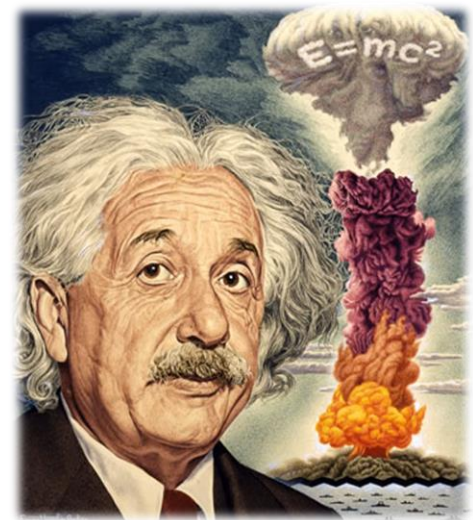
54 mm



**L = 51 mm @ 20  $\mu\text{m}$**

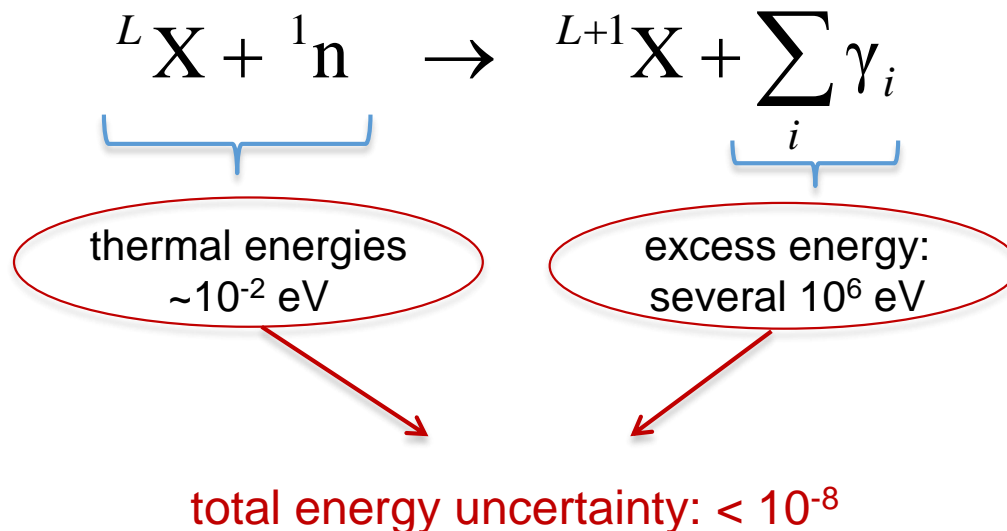


$$E = mc^2$$



How can we test it?

- need process, where mass is converted in energy
- thermal-neutron capture reaction:





$$[m(n) + m({}^L\text{X}) - m({}^{L+1}\text{X})]c^2 = \sum_i E(\gamma_i) \quad ?$$

$$E(\gamma) = h\nu = \frac{hc}{\lambda}$$

In terms of mass units  $A$  relative to an atomic mass scale  $u = 10^{-3}/N_A$  kg :

$$A(n) + \underbrace{A({}^L\text{X}) - A({}^{L+1}\text{X})}_{\Delta A({}^{L+1}\text{X})} = 10^3 \frac{N_A h}{c} \sum_i \frac{1}{\lambda_i({}^{L+1}\text{X})}$$

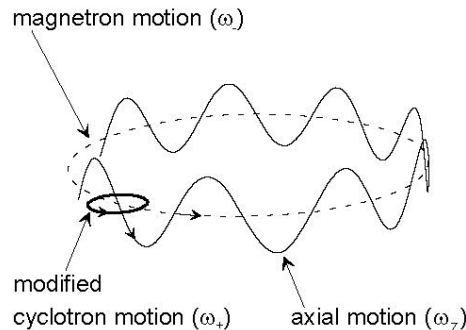
molar Planck constant

$$\underbrace{\Delta A({}^{L+1}\text{X}) - \Delta A({}^{K+1}\text{Y})}_{\text{Penning trap measurements (4 masses)}} = 10^3 \frac{N_A h}{c} \left[ \sum_i \frac{1}{\lambda_i({}^{L+1}\text{X})} - \sum_j \frac{1}{\lambda_j({}^{K+1}\text{Y})} \right]$$

Gamma-ray wavelength measurements  
(2 nuclides after neutron capture)



$$\omega = \frac{qB}{m}$$

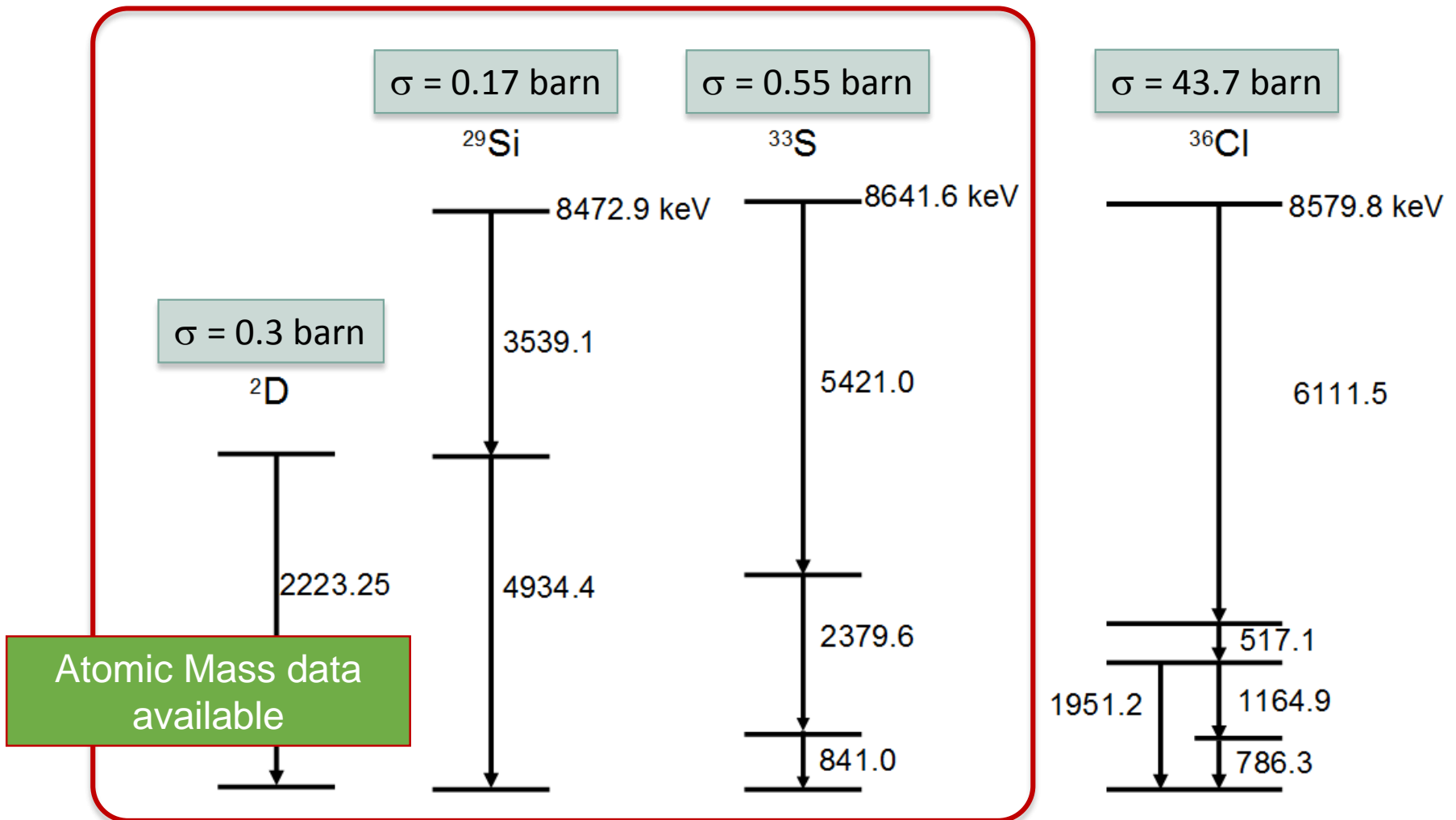


Double crystal monochromator  
**GAMS**

# Which isotopes?

Penning Trap:  $A(L, L+1X)$  can be measured with  $10^{-11}$  relative uncertainty!

Need mass values for two pairs of stable isotopes



A  $10^{-11}$  relative uncertainty on  $A$  requires  $\lambda$  measurements with accuracy  $10^{-8}$

What is needed to determine gamma energies (or wavelengths) with high accuracy?



is necessary but insufficient

High accuracy  $\rightarrow$  use gamma spectroscopy based on **Laue diffraction**

$$n \frac{hc}{E_\gamma} = 2d \sin \theta$$

Bragg's law for photons

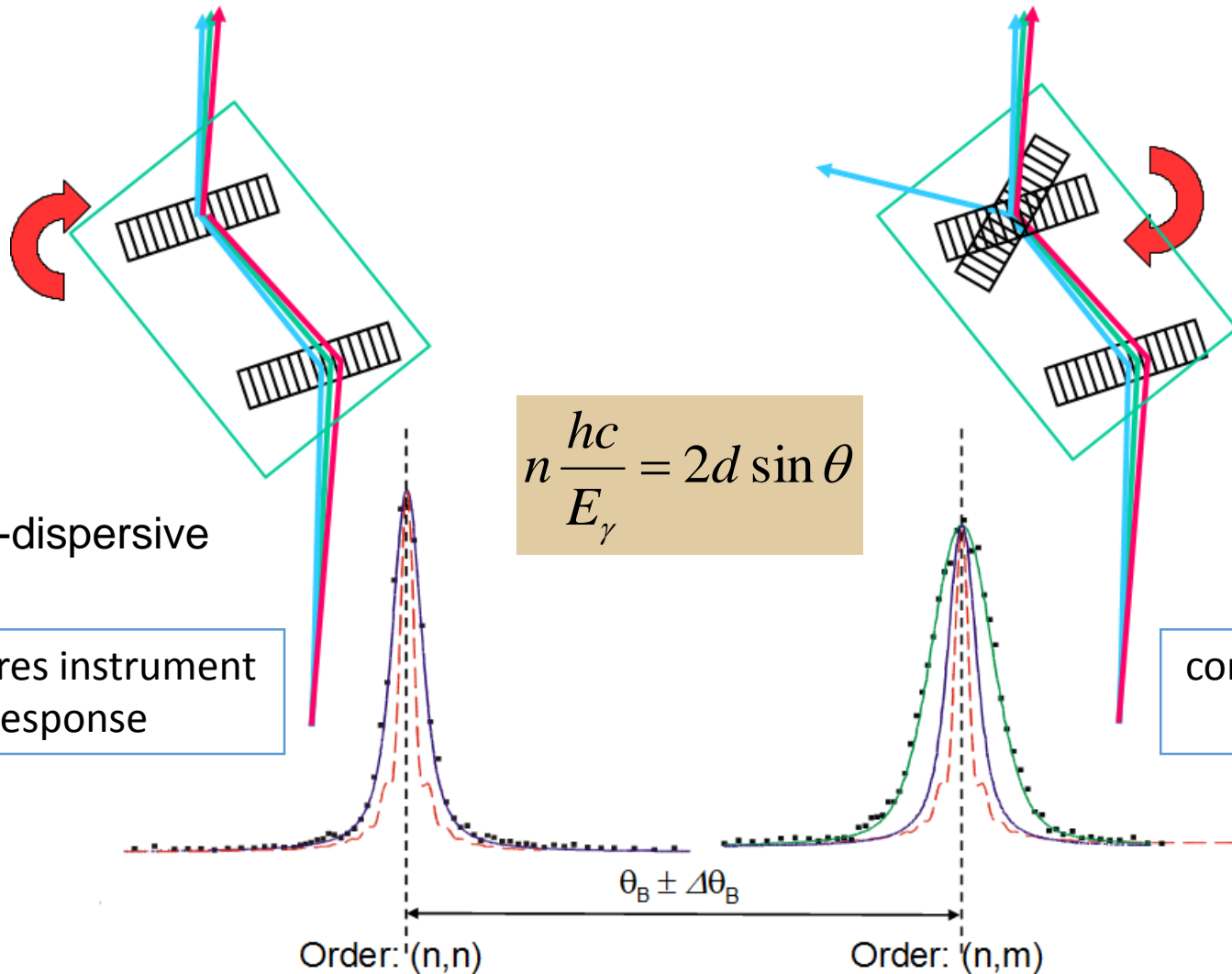
Need absolute measurements of:

- lattice constant  $d$
- scattering angle  $\theta$

$$\left( \frac{\Delta E_\gamma}{E_\gamma} \right)^2 \sim \left( \frac{\Delta \theta}{\theta} \right)^2 + \left( \frac{\Delta d}{d} \right)^2$$

*$\Delta d/d < 10^{-8}$*

# Flat double-crystal monochromator spectrometer for gamma rays: **GAMS**



non-dispersive

measures instrument  
response

$$n \frac{hc}{E_\gamma} = 2d \sin \theta$$

dispersive

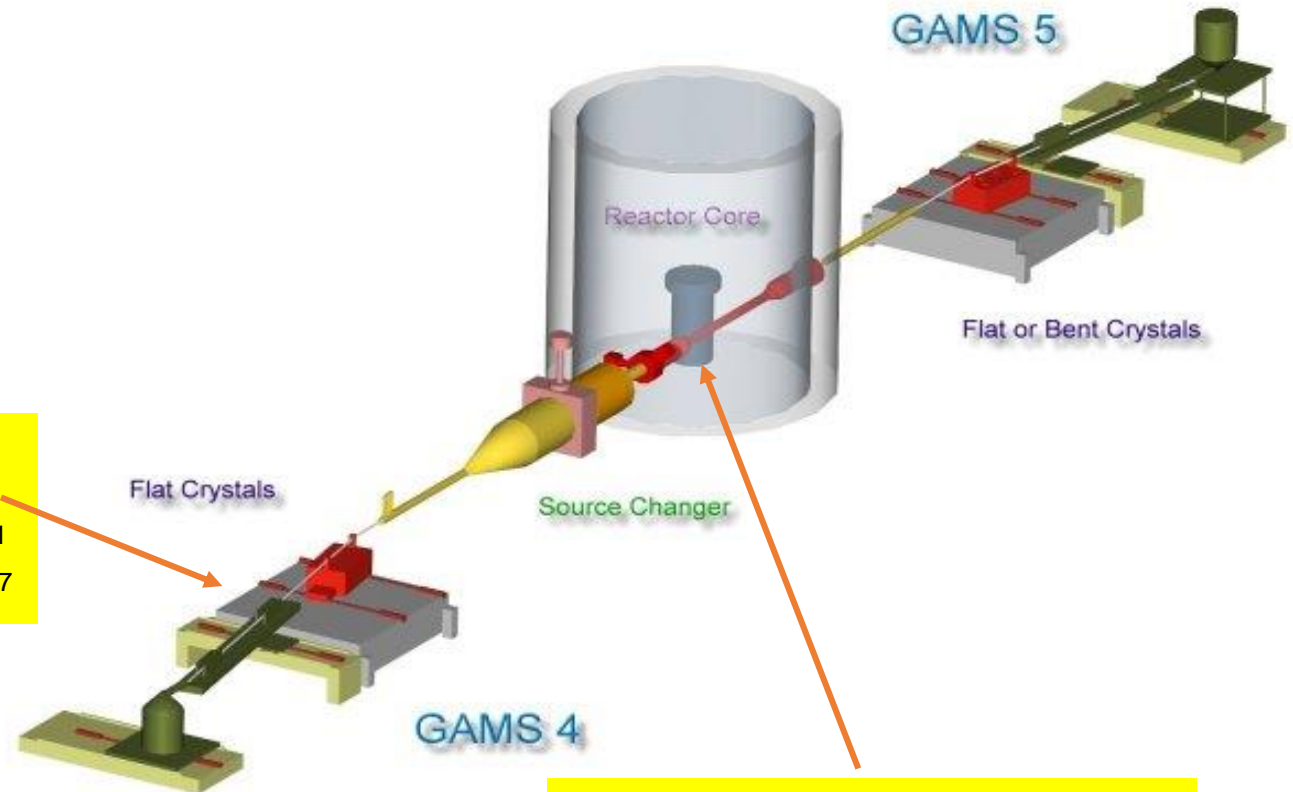
contains additional  
broadening

$\theta_B \pm \Delta\theta_B$

Order: (n,n)

Order: (n,m)

# Implantation of spectrometer

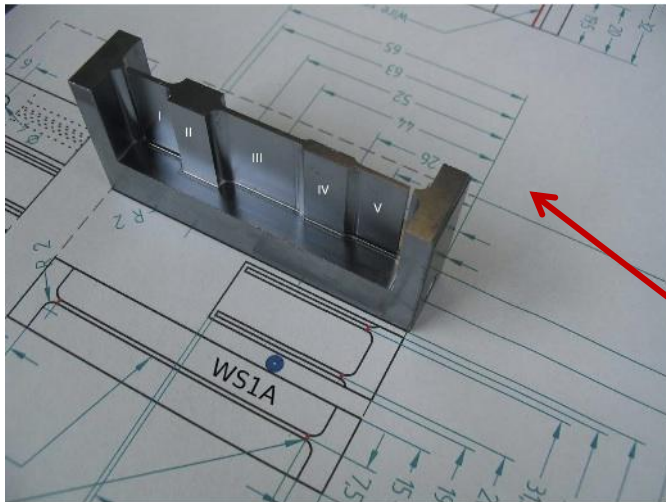
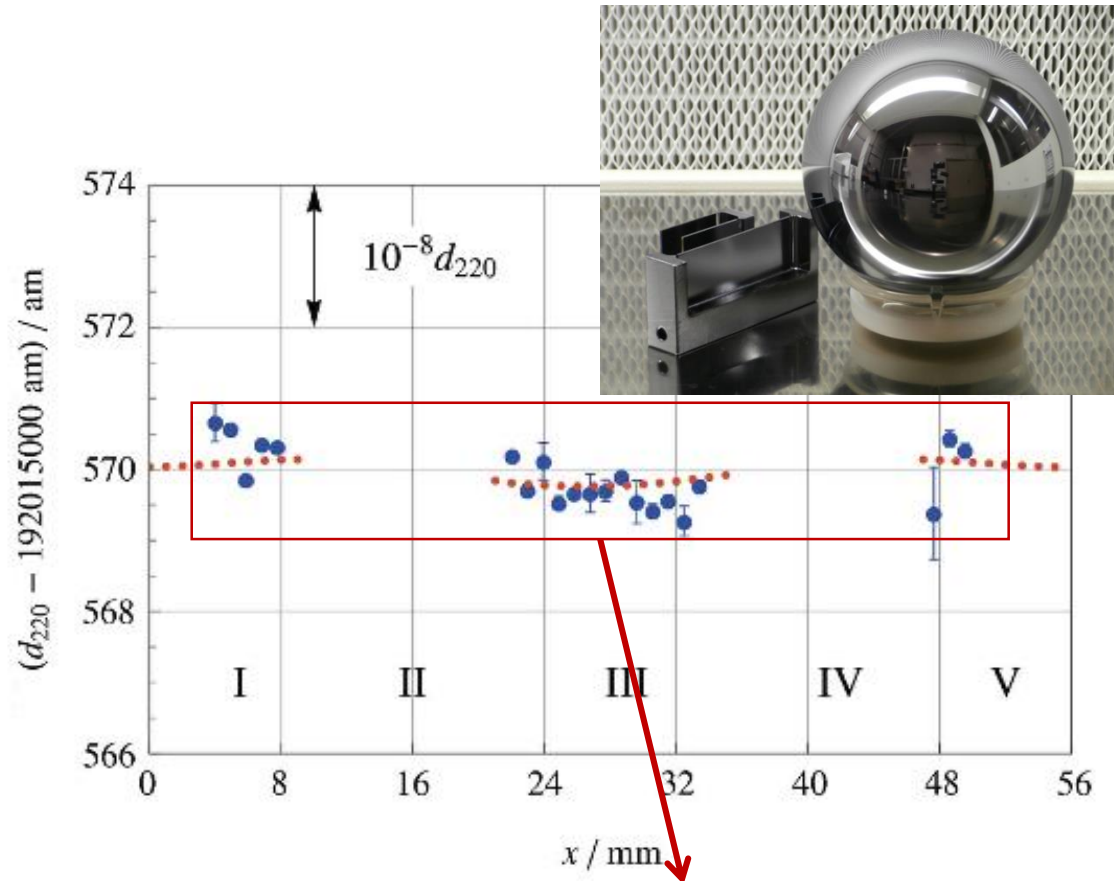
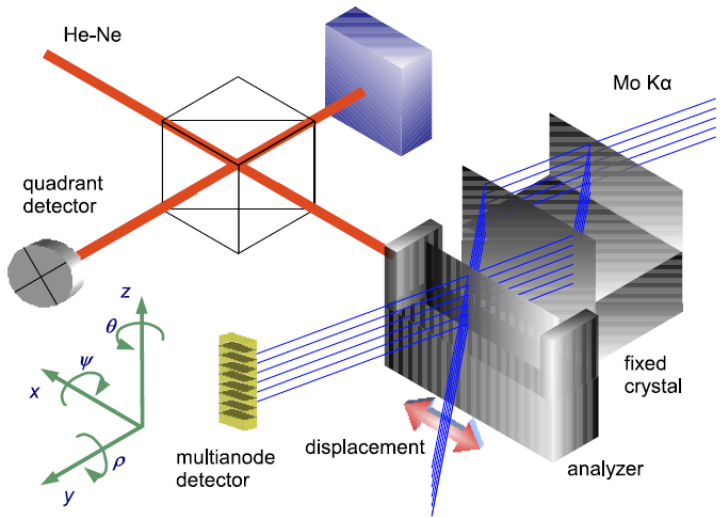


## Flat Crystals:

- Resolution:  $10^{-6}$
- Eff. Solid Angle:  $10^{-11}$
- absolute Energy:  $10^{-7}$

Neutron Flux:  $5 \times 10^{14}$   
Targets: 0.1 – 10g  
Target change during reactor cycle

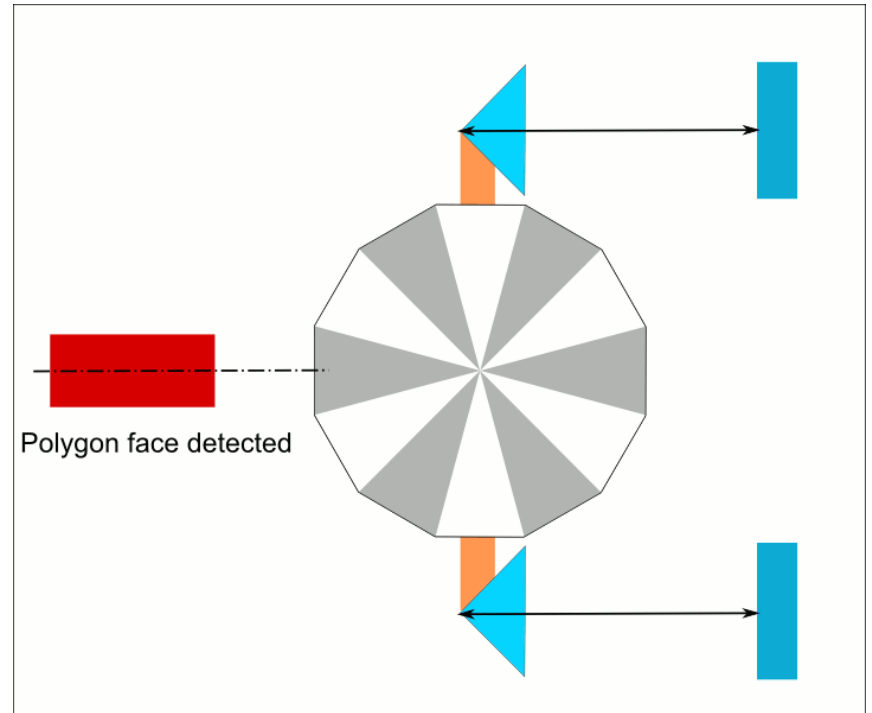
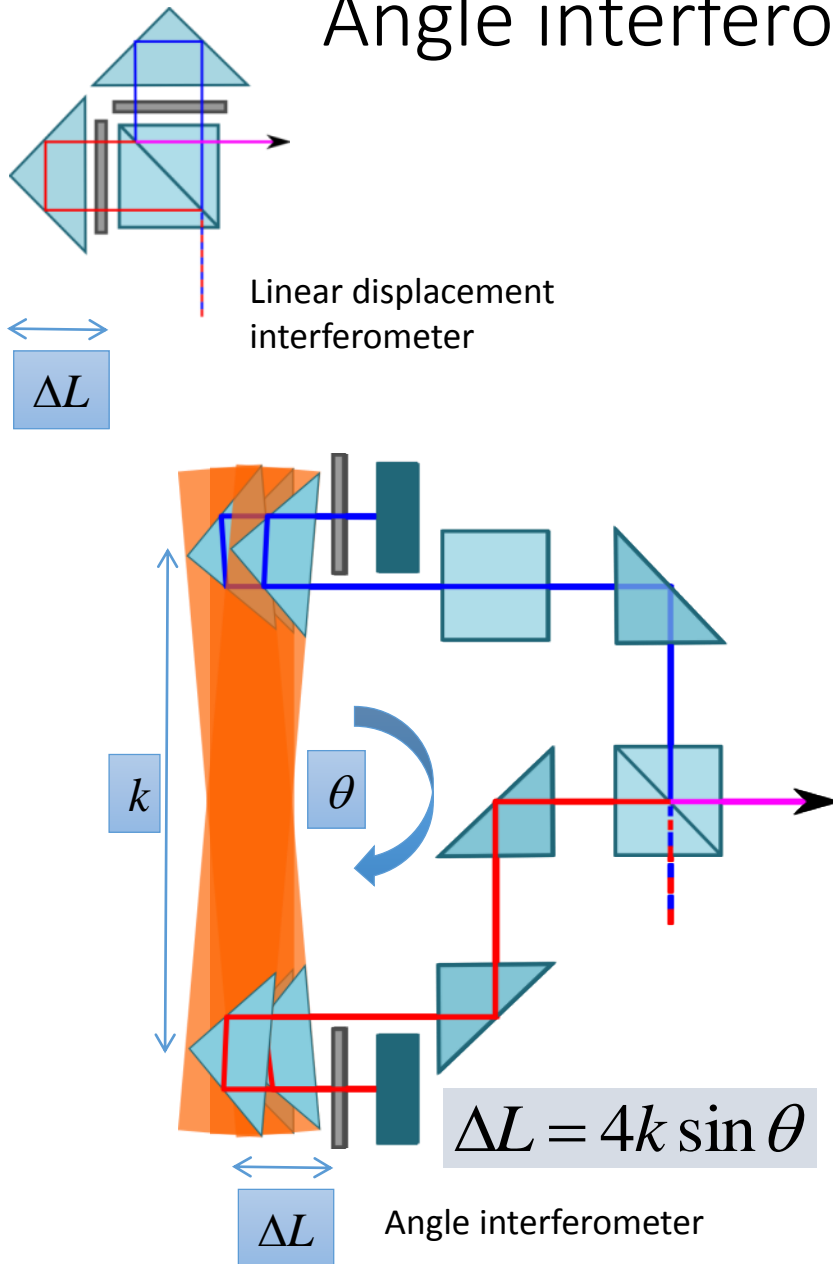
# $\Delta d/d$ : How perfect are GAMS crystals?



$\Delta d/d < 10^{-8}$ , extremely perfect silicon crystals,  $d$  known in SI units

- Silicon crystal for GAMS6 from WASO4 reference material
- Fabricated by PTB, Braunschweig, Germany
  - Characterized by INRIM, Torino, Italy

# Angle interferometer with self-calibration



$$2\pi = \sum_i^N \arcsin\left(\frac{\Delta L_i}{4k}\right)$$

- $N \Delta L_i$  are measured
- Equation is solved for  $k$

# Result of $E = mc^2$ test using GAMS4

$$E_\gamma - \Delta mc^2 = -(1.2 \pm 4.3) \times 10^{-7}$$

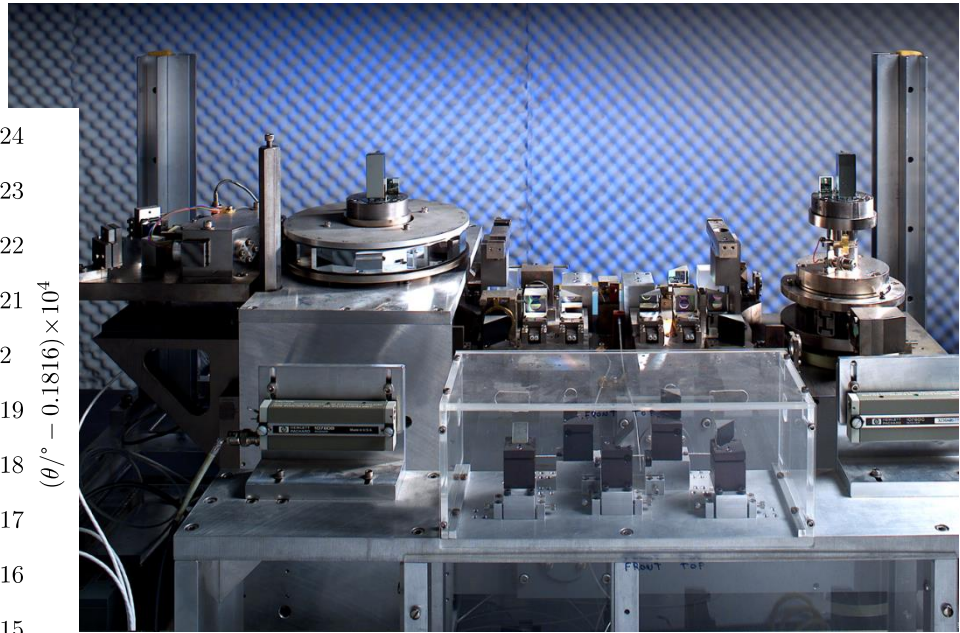
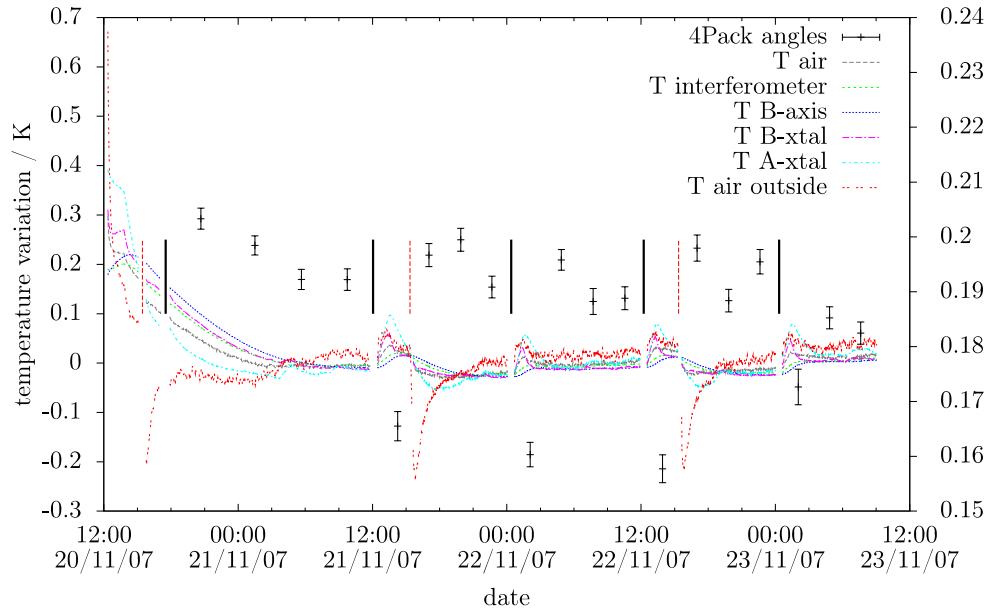
S. Rainville et al., [Nature](#) **438** (2005) 1096

$A+1X$	$\Delta m$ from Penning trap (u)	$\Delta m$ from GAMS4 (u)	Rel. Diff. $\times 10^7$
$^{29}\text{Si}$	0.00670861569(47)	0.00670860929(536)	-9.5(8.0)
$^{33}\text{S}$	0.00688901053(50)	0.00688901206(351)	2.2(5.1)
weighted average relative difference			-1.2(4.3)

Wavelength measurements are limiting



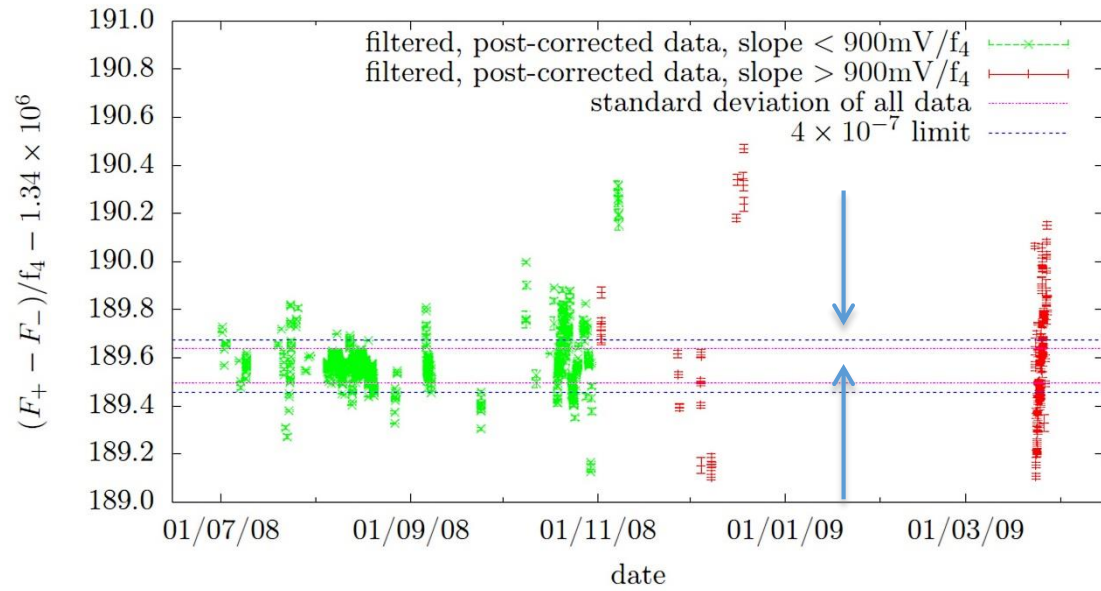
# Accuracy reach of GAMS4



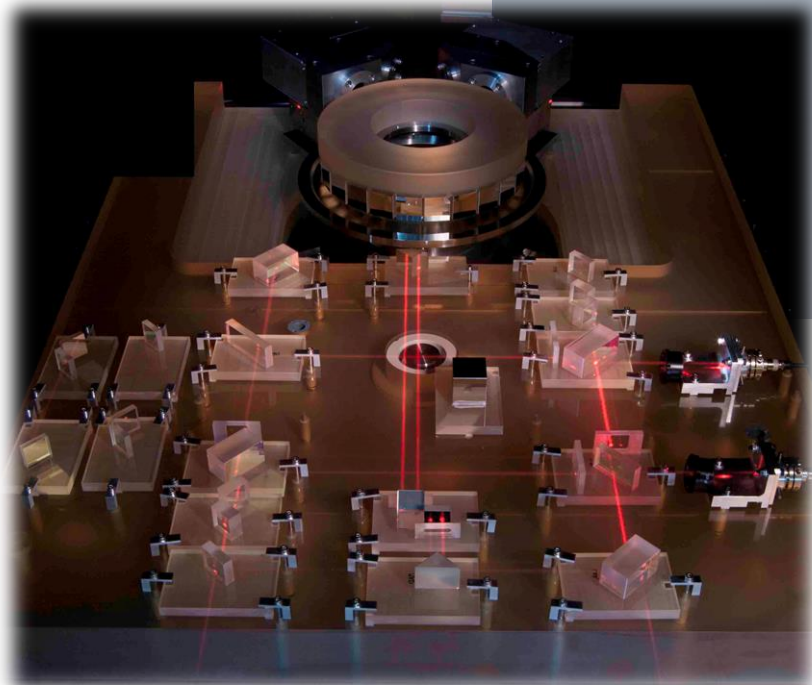
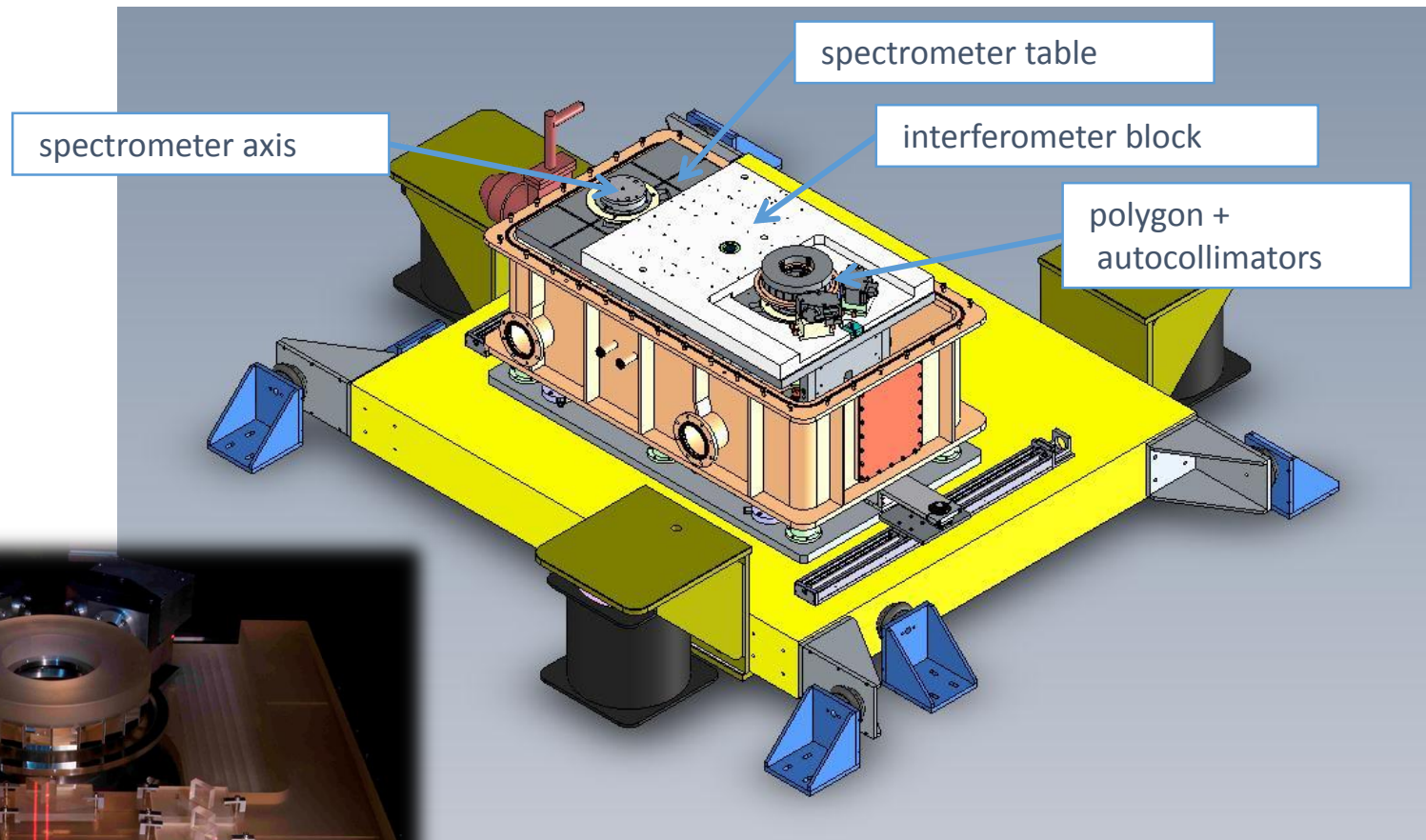
Stability of calibration:  $2.1 \times 10^{-7}$

Angle measurement and calibration under atmosphere possible not better than on  $10^{-7}$  level

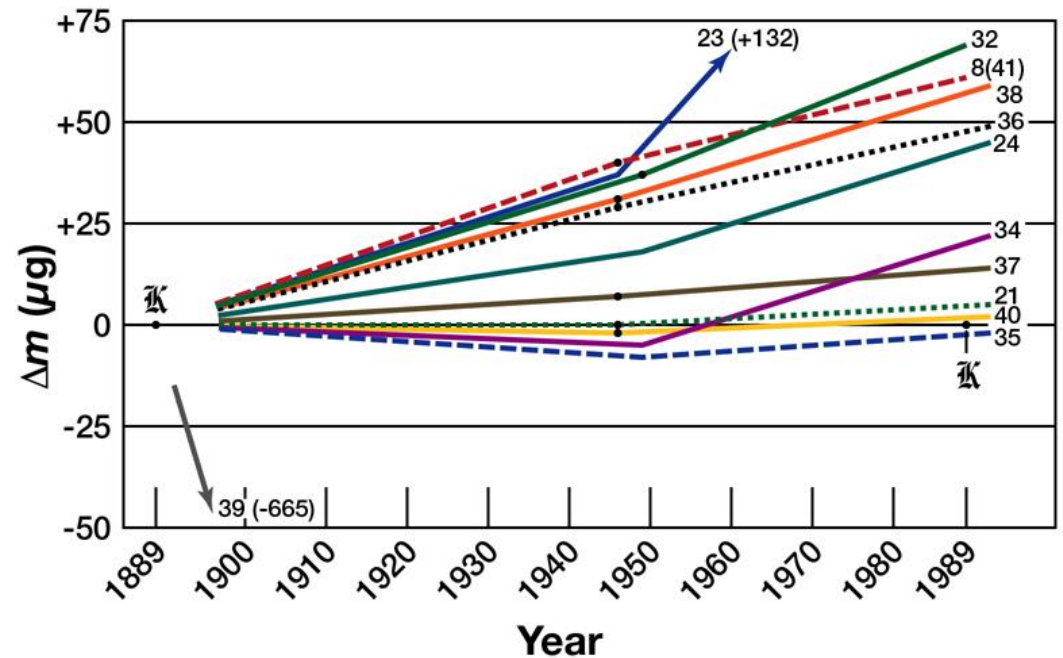
J. Krempel, PhD LMU, 2011



# New instrument GAMS6

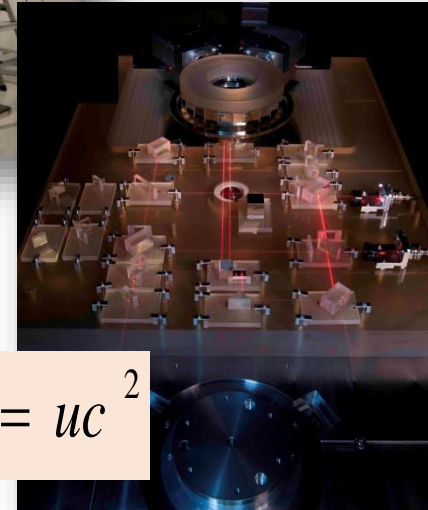
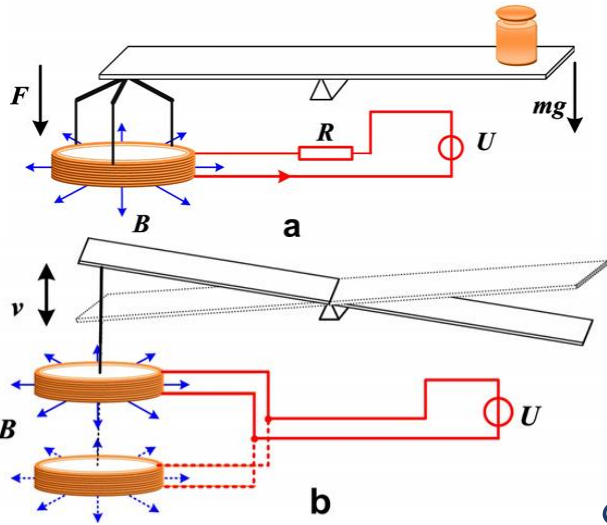


# Challenge: redefinition of the kilogram



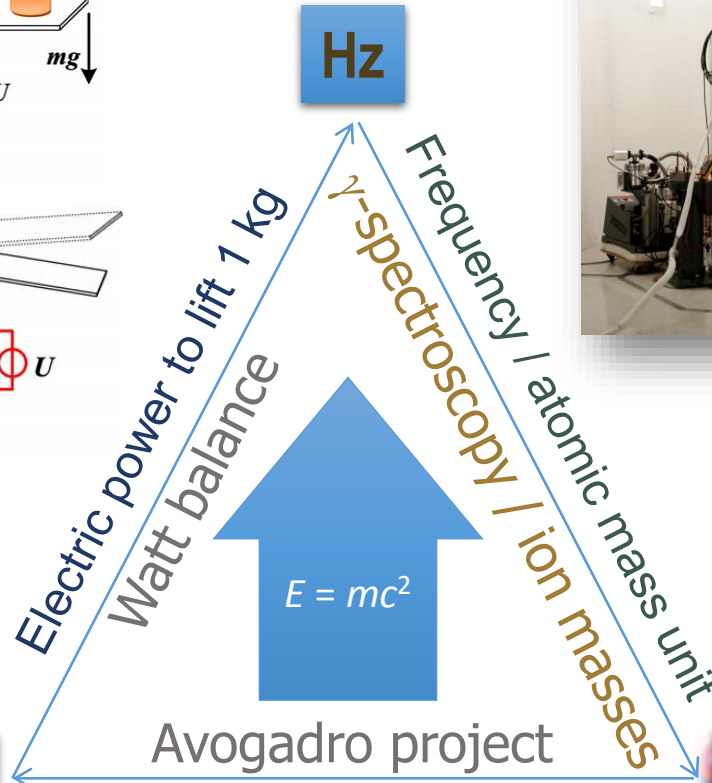


# Routes to a new mass unit definition



$$h = \frac{4m_{\text{kg}} g v}{R_{J-90}^2 R_{K-90} (UI)_{90}}$$

$$N_A h \nu = u c^2$$



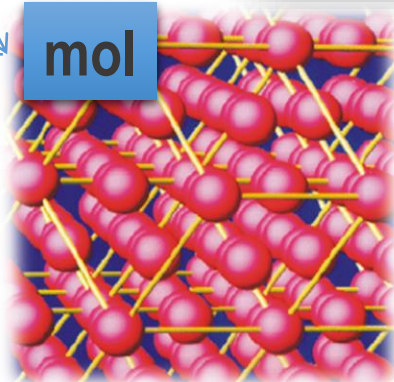
$$E = mc^2$$



**kg**

Number of atoms in 1 kg

$$m_{\text{kg}} = 10^3 \times N_A u$$



**mol**

# Acknowledgements

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*Technische Universität München*

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**Maurits van der Grinten**  
*Rutherford Appleton Laboratory*



## Short-range gravity

**Hartmut Abele**  
*Atominstytut Wien*



$E = mc^2$   
**Michael Jentschel**  
*Institut Laue Langevin*